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The Hilbert space costratification for the orbit type strata of $SU(2)$ -lattice gauge theory

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We construct the Hilbert space costratification of $G = SU(2)$ -quantum gauge theory on a finite spatial lattice in the Hamiltonian approach. We build on previous work [F. Fürstenberg, G. Rudolph, and M. Schmidt, *J. Geom. Phys.* **119**, 66–81 (2017)], where we have implemented the classical gauge orbit strata on the quantum level within a suitable holomorphic picture. In this picture, each element τ of the classical stratification corresponds to the zero locus of a finite subset $\{p_i\}$ of the algebra \mathcal{R} of G -invariant representative functions on $G_{\mathbb{C}}^N$. Viewing the invariants as multiplication operators \hat{p}_i on the Hilbert space \mathcal{H} , the union of their images defines a subspace of \mathcal{H} whose orthogonal complement \mathcal{H}_τ is the element of the costratification corresponding to τ . To construct \mathcal{H}_τ , one has to determine the images of the \hat{p}_i explicitly. To accomplish this goal, we construct an orthonormal basis in \mathcal{H} and determine the multiplication law for the basis elements; that is, we determine the structure constants of \mathcal{R} in this basis. This part of our analysis applies to any compact Lie group G . For $G = SU(2)$, the above procedure boils down to a problem in combinatorics of angular momentum theory. Using this theory, we obtain the union of the images of the operators \hat{p}_i as a subspace generated by vectors whose coefficients with respect to our basis are given in terms of Wigner's $3nj$ symbols. The latter are further expressed in terms of $9j$ symbols. Using these techniques, we are also able to reduce the eigenvalue problem for the Hamiltonian of this theory to a problem in linear algebra. *Published by AIP Publishing.* <https://doi.org/10.1063/1.5031115>

I. INTRODUCTION

This paper is a continuation of our previous one,¹¹ where we have derived the defining relations for the orbit type strata of $G = SU(2)$ -lattice gauge theory. It is part of a program which aims at developing a non-perturbative approach to the quantum theory of gauge fields in the Hamiltonian framework with special emphasis on the role of non-generic gauge orbit types. The starting point is a finite-dimensional Hamiltonian lattice approximation of the theory, which on the classical level leads to a finite-dimensional Hamiltonian system with symmetries. The corresponding quantum theory is obtained via canonical quantization. It is best described in the language of C^* -algebras with a field algebra which (for a pure gauge theory) may be identified with the algebra of compact operators on the Hilbert space of square-integrable functions over the product G^N of N copies of the gauge group manifold G , where N is the number of off-tree links for a chosen maximal lattice tree. Correspondingly, the observable algebra is obtained via gauge symmetry reduction. We refer to Refs. 23, 25, 26, and 34 for the study of this algebra, including its superselection structure. For first steps toward the construction of the thermodynamical limit, see Refs. 13 and 14.

If the gauge group is non-Abelian, then the action of the symmetry group in the corresponding classical Hamiltonian system necessarily has more than one orbit type. Correspondingly, the reduced phase space obtained by symplectic reduction is a stratified symplectic space^{32,34,36} rather than a symplectic manifold as in the case with one orbit type.¹ The stratification is given by the orbit type

strata. It consists of an open and dense principal stratum and several secondary strata. Each of these strata is invariant under the dynamics with respect to any invariant Hamiltonian. For case studies, we refer to Refs. 6, 7, and 9.

To study the influence of the classical orbit type stratification on the quantum level, we use the concept of costratification of the quantum Hilbert space as developed by Huebschmann.²⁰ A costratification is given by a family of closed subspaces, one for each stratum. Loosely speaking, the closed subspace associated with a certain classical stratum consists of the wave functions which are optimally localized at that stratum in the sense that they are orthogonal to all states vanishing at that stratum. The vanishing condition can be given sense in the framework of holomorphic quantization, where wave functions are true functions and not just classes of functions. In Ref. 22, we have constructed this costratification for a toy model with gauge group $SU(2)$ on a single lattice plaquette. As physical effects, we have found a nontrivial overlap between distant strata and, for a certain range of the coupling, a very large transition probability between the ground state of the lattice Hamiltonian and one of the two secondary strata.

In the present paper, we deal with the theory with gauge group $G = SU(2)$ for an arbitrary finite lattice. In this case, there are non-trivial relations characterizing the classical gauge orbit strata which, in a first step, should be implemented on the quantum level. This problem has been solved in Ref. 11 using the above mentioned holomorphic picture. In this picture, each element τ of the stratification corresponds to the zero locus of a finite subset $\{p_1, \dots, p_r\}$ of the algebra \mathcal{R} of G -invariant representative functions on $G_{\mathbb{C}}^N$, where $G_{\mathbb{C}}$ denotes the complexification of G . Viewing the invariants p_i as multiplication operators \hat{p}_i on the Hilbert space \mathcal{H} , the union of their images defines a subspace of \mathcal{H} whose orthogonal complement \mathcal{H}_{τ} is, by definition, the element of the costratification corresponding to τ . Thus, to construct \mathcal{H}_{τ} , one has to determine the images of the \hat{p}_i explicitly. To accomplish this goal, we construct an orthonormal basis in \mathcal{H} and determine the structure constants of the algebra \mathcal{R} with respect to that basis. This part of our analysis applies to any compact Lie group G . So, assuming that we are given the classical stratification for some Lie group G in terms of the classical invariants p_i , with the above result at our disposal, we can in principle determine the operators \hat{p}_i as well as their images in \mathcal{H} in terms of linear combinations of the elements of the chosen basis.

For $G = SU(2)$, our procedure boils down to a problem in combinatorics of angular momentum theory. For the latter, we refer to Refs. 4, 5, 28, 37, and 41. Using this theory, we obtain the union of the images of the operators \hat{p}_i as a subspace of \mathcal{H} generated by vectors whose coefficients with respect to our basis are given in terms of Wigner's $3nj$ symbols. The latter is further expressed in terms of $9j$ symbols. For these symbols, there exist nowadays efficient calculators; that is, the above coefficients can be calculated explicitly. Using the same techniques, we are also able to reduce the eigenvalue problem for the Hamiltonian of this theory to a problem in linear algebra.

The paper is organized as follows. In Sec. II, we explain the model. To keep the presentation self-contained, in Sec. III, we present the basics of stratified quantum gauge theory for arbitrary compact gauge groups as developed in detail in Ref. 11. Moreover, we construct a basis of \mathcal{H} consisting of G -invariant representative functions and derive the multiplication law in \mathcal{R} . In Sec. IV, we turn to the study of the case $G = SU(2)$. We describe the orbit type strata, analyze the zero loci in terms of the above basis, and obtain the images of the multiplication operators \hat{p}_i in terms of linear combinations of the basis elements. As a consequence, the costrata are given by systems of linear equations with real coefficients built from $3nj$ symbols. We illustrate the result for the special case $N = 2$. Finally, we discuss the eigenvalue problem for the Hamiltonian in terms of the above basis. The text is completed by Appendixes A and B containing the proofs of two technical results.

II. THE MODEL

Let G be a compact Lie group and let \mathfrak{g} be its Lie algebra. Later on, we will specify $G = SU(2)$, but for the time being, this is not necessary. Let Λ be a finite spatial lattice and let Λ^0 , Λ^1 , and Λ^2 denote, respectively, the sets of lattice sites, lattice links, and lattice plaquettes. For the links and plaquettes, let there be chosen an arbitrary orientation. In lattice gauge theory with gauge group G in the Hamiltonian approach, gauge fields (the variables) are approximated by their parallel transporters

along links and gauge transformations (the symmetries) are approximated by their values at the lattice sites. Thus, the classical configuration space is the space G^{Λ^1} of mappings $\Lambda^1 \rightarrow G$, the classical symmetry group is the group G^{Λ^0} of mappings $\Lambda^0 \rightarrow G$ with pointwise multiplication, and the action of $g \in G^{\Lambda^0}$ on $a \in G^{\Lambda^1}$ is given by

$$(g \cdot a)(\ell) := g(x)a(\ell)g(y)^{-1}, \quad (2.1)$$

where $\ell \in \Lambda^1$ and x, y denote the starting point and the endpoint of ℓ , respectively. The classical phase space is given by the associated Hamiltonian G -manifold,^{1,35} and the reduced classical phase space is obtained from that by symplectic reduction.^{32,35,36} We do not need the details here. Dynamics is ruled by the classical counterpart of the Kogut-Susskind lattice Hamiltonian; see Subsection IV F. When identifying T^*G with $G \times \mathfrak{g}$, and thus $T^*G^{\Lambda^1}$ with $G^{\Lambda^1} \times \mathfrak{g}^{\Lambda^1}$, by means of left-invariant vector fields, the classical Hamiltonian is given by

$$H(a, E) = \frac{g^2}{2\delta} \sum_{\ell \in \Lambda^1}^N \|E(\ell)\|^2 - \frac{1}{g^2\delta} \sum_{p \in \Lambda^2} (\text{tr } a(p) + \overline{\text{tr } a(p)}), \quad (2.2)$$

where $a \in G^{\Lambda^1}$, g denotes the coupling constant, δ denotes the lattice spacing, and $a(p)$ denotes the product of $a(\ell)$ along the boundary of the plaquette p in the induced orientation. The trace is taken in some chosen unitary representation. Unitarity ensures that the Hamiltonian does not depend on the choice of plaquette orientations. Finally, $E \in \mathfrak{g}^{\Lambda^1}$ is the classical colour electric field (canonically conjugate momentum).

When discussing orbit types in continuum gauge theory, it is convenient to first factorize with respect to the free action of pointed gauge transformations, thus arriving at an action of the compact gauge group G on the quotient manifold. This preliminary reduction can also be carried out in the case of lattice gauge theory under consideration. In fact, given a lattice site x_0 , it is not hard to see that the normal subgroup

$$\{g \in G^{\Lambda^0} : g(x_0) = \mathbb{1}\}, \quad (2.3)$$

where $\mathbb{1}$ denotes the unit element of G , acts freely on G^{Λ^1} . Hence, one may pass to the quotient manifold and the residual action by the quotient Lie group of G^{Λ^0} with respect to this normal subgroup. Clearly, the quotient Lie group is naturally isomorphic to G . The quotient manifold can be identified with a direct product of copies of G , and the quotient action can be identified with the action of G by diagonal conjugation as follows. Choose a maximal tree \mathcal{T} in the graph Λ^1 and define the tree gauge of \mathcal{T} to be the subset

$$\{a \in G^{\Lambda^1} : a(\ell) = \mathbb{1} \ \forall \ell \in \mathcal{T}\}$$

of G^{Λ^1} . One can readily see that every element of G^{Λ^1} is conjugate under G^{Λ^0} to an element in the tree gauge of \mathcal{T} and that two elements in the tree gauge of \mathcal{T} are conjugate under G^{Λ^0} if they are conjugate under the action of G via constant gauge transformations. This implies that the natural inclusion mapping of the tree gauge into G^{Λ^1} descends to a G -equivariant diffeomorphism from that tree gauge onto the quotient manifold of G^{Λ^1} with respect to the action of the subgroup (2.3). Finally, by choosing a numbering of the off-tree links in Λ^1 , we can identify the tree gauge of \mathcal{T} with the direct product of N copies of G , where N denotes the number of off-tree links. This number does not depend on the choice of \mathcal{T} . Then, the action of G on the tree gauge via constant gauge transformations translates into the action of G on G^N by diagonal conjugation,

$$g \cdot (a_1, \dots, a_N) = (ga_1g^{-1}, \dots, ga_Ng^{-1}). \quad (2.4)$$

As a consequence of these considerations, for the discussion of the role of orbit types, we may pass from the original large Hamiltonian system with symmetries, given by the configuration space G^{Λ^1} , the symmetry group G^{Λ^0} , and the action (2.1), to the smaller Hamiltonian system with symmetries given by the configuration space

$$Q := G^N,$$

the symmetry group G , and the action of G on Q given by diagonal conjugation (2.4). This is the system we will discuss here. As before, the classical phase space is given by the associated Hamiltonian

G -manifold and the reduced classical phase space is obtained from that by symplectic reduction. One can show that the latter is isomorphic, as a stratified symplectic space, to the reduced classical phase space defined by the original Hamiltonian system with symmetries.

We will need the following information about the classical phase space. As a space, it is given by the cotangent bundle

$$T^*Q \equiv T^*G^N.$$

It is a general fact that the action of G on Q naturally lifts to a symplectic action on T^*Q (consisting of the corresponding ‘point transformations’ in the language of canonical transformations) and that the lifted action admits a momentum mapping

$$\mu : T^*Q \rightarrow \mathfrak{g}^*, \quad \mu(p)(A) := p(A_*),$$

where $p \in T^*Q$, $A \in \mathfrak{g}$, and A_* denotes the Killing vector field defined by A . An easy calculation shows that under the global trivialization

$$T^*G^N \cong G^N \times \mathfrak{g}^N \quad (2.5)$$

induced by left-invariant vector fields and an invariant scalar product on \mathfrak{g} , the lifted action is given by diagonal conjugation,

$$g \cdot (a_1, \dots, a_N, A_1, \dots, A_N) = (ga_1g^{-1}, \dots, ga_Ng^{-1}, \text{Ad}(g)A_1, \dots, \text{Ad}(g)A_N), \quad (2.6)$$

and the associated momentum mapping is given by

$$\mu(a_1, \dots, a_N, A_1, \dots, A_N) = \sum_{i=1}^N \text{Ad}(a_i)A_i - A_i, \quad (2.7)$$

see, e.g., Ref. 35, Sec. 10.7. The reduced phase space \mathcal{P} is obtained from T^*G^N by singular symplectic reduction at $\mu = 0$. That is, \mathcal{P} is the set of orbits of the lifted action of G on the invariant subset $\mu^{-1}(0) \subset T^*Q$, endowed with the quotient topology induced from the relative topology on this subset. In lattice gauge theory, the condition $\mu = 0$ corresponds to the Gauß law constraint. As a matter of fact, the action of G on $\mu^{-1}(0)$ has the same orbit types as that on Q . By definition, the orbit type strata of \mathcal{P} are the connected components of the subsets of \mathcal{P} of elements with a fixed orbit type. They are called strata because they provide a stratification of \mathcal{P} .^{32,36} By the procedure of symplectic reduction, the orbit type strata of \mathcal{P} are endowed with symplectic manifold structures. The bundle projection $T^*Q \rightarrow Q$ induces a mapping $\mathcal{P} \rightarrow Q/G$. This mapping is surjective because μ is linear on the fibres of T^*Q and hence $\mu^{-1}(0)$ contains the zero section of T^*Q . It need not preserve the orbit type though.

Remark 2.1. The tree gauge need not be invariant under time evolution with respect to a gauge-invariant Hamiltonian (e.g., the Kogut-Susskind lattice Hamiltonian), but every motion in the full configuration space G^{Λ^1} can be transformed by a time-dependent gauge transformation to the tree gauge. Thus, up to time-dependent gauge transformations, the tree gauge is invariant under time evolution. This is reflected in the isomorphism of the reduced phase spaces mentioned above. ♦

III. STRATIFIED QUANTUM THEORY

A. Quantization and reduction

To construct the quantum theory of the reduced system, one may either first reduce the classical system and then quantize or first quantize and then reduce the quantum system. Here, we follow the second strategy; that is, we carry out geometric (Kähler) quantization on T^*G^N and subsequent reduction. Let $\mathfrak{g}_{\mathbb{C}}$ denote the complexification of \mathfrak{g} and let $G_{\mathbb{C}}$ denote the complexification of G . This is a complex Lie group having G as its maximal compact subgroup. It is unique up to isomorphism. For $G = \text{SU}(n)$, we have $G_{\mathbb{C}} = \text{SL}(n, \mathbb{C})$. By restriction, the exponential mapping

$$\exp: \mathfrak{g}_{\mathbb{C}} \rightarrow G_{\mathbb{C}}$$

of $G_{\mathbb{C}}$ and multiplication in $G_{\mathbb{C}}$ induce a diffeomorphism

$$G \times \mathfrak{g} \rightarrow G^{\mathbb{C}}, \quad (a, A) \mapsto a \exp(iA), \quad (3.1)$$

which is equivariant with respect to the action of G on $G \times \mathfrak{g}$ by

$$g \cdot (a, A) := (gag^{-1}, \text{Ad}(g)A)$$

and the action of G on $G_{\mathbb{C}}$ by conjugation. For $G = \text{SU}(n)$, this diffeomorphism amounts to the inverse of the polar decomposition. By applying this diffeomorphism to each copy, we obtain a diffeomorphism

$$G^N \times \mathfrak{g}^N \rightarrow G_{\mathbb{C}}^N, \quad (a_1, \dots, a_N, A_1, \dots, A_N) \mapsto (a_1 \exp(iA_1), \dots, a_N \exp(iA_N)).$$

By composing the latter with the global trivialization (2.5), we obtain a diffeomorphism

$$T^*G^N \rightarrow G_{\mathbb{C}}^N \quad (3.2)$$

which, due to (2.6), is equivariant with respect to the lifted action of G on T^*G^N and the action of G on $G_{\mathbb{C}}^N$ by diagonal conjugation. Via this diffeomorphism, the complex structure of $G_{\mathbb{C}}^N$ and the symplectic structure of T^*G^N combine to a Kähler structure. Half-form Kähler quantization on $G_{\mathbb{C}}^N$ yields the Hilbert space

$$HL^2(G_{\mathbb{C}}^N, d\nu_{\hbar})$$

of holomorphic functions on $G_{\mathbb{C}}^N$ which are square-integrable with respect to the measure

$$d\nu_{\hbar} = e^{-\kappa/\hbar} \eta \varepsilon, \quad (3.3)$$

where

$$\kappa(a_1 e^{iA_1}, \dots, a_N e^{iA_N}) = |A_1|^2 + \dots + |A_N|^2$$

is the Kähler potential on $G_{\mathbb{C}}^N$,

$$\eta(a_1 e^{iA_1}, \dots, a_N e^{iA_N}) = \sqrt{\det \frac{\sin(\text{ad}(A_1))}{\text{ad}(A_1)}} \dots \sqrt{\det \frac{\sin(\text{ad}(A_N))}{\text{ad}(A_N)}}$$

is the half-form correction, and

$$\varepsilon(a_1 e^{iA_1}, \dots, a_N e^{iA_N}) = da_1 \dots da_N dA_1 \dots dA_N$$

is the Liouville measure on T^*G^N . Reduction then yields the closed subspace

$$\mathcal{H} = HL^2(G_{\mathbb{C}}^N, d\nu_{\hbar})^G$$

of G -invariants as the Hilbert space of the reduced system.

Remark 3.1. The above result belongs to Hall.¹⁶ Alternatively, the Hilbert space $HL^2(G_{\mathbb{C}}^N, d\nu)$ is obtained via the Segal-Bargmann transformation for compact Lie groups.¹⁵ In more detail, the Segal-Bargmann transformation

$$\Phi: L^2(G^N) \rightarrow HL^2(G_{\mathbb{C}}^N, d\nu_{\hbar})$$

is a unitary isomorphism, which restricts to a unitary isomorphism of the subspaces of invariants. ♦

B. Orbit type costratification

Following Huebschmann,²⁰ we define the subspaces associated with the orbit type strata of \mathcal{P} to be the orthogonal complements of the subspaces of functions vanishing on those strata. To accomplish this idea, we first clarify how to interpret elements of \mathcal{H} as functions on \mathcal{P} . In the case $N = 1$ discussed in Refs. 22 and 19, this is readily done by observing that $\mathcal{P} \cong T_{\mathbb{C}}/W$, where T is a maximal

torus in G and W is the corresponding Weyl group, and using the isomorphism $HL^2(G_{\mathbb{C}}, d\nu)^G \cong HL^2(T_{\mathbb{C}}, d\nu_T)^W$; see Sec. 3.1 of Ref. 22. Here, the measure $d\nu_T$ is obtained from $d\nu$ by integration over the conjugation orbits in $G_{\mathbb{C}}$, thus yielding an analog of Weyl's integration formula for $HL^2(G_{\mathbb{C}}, d\nu)$. In the general case, the argument is as follows.

First, we construct a quotient of $G_{\mathbb{C}}^N$ on which the elements of \mathcal{H} define functions. Consider the action of $G_{\mathbb{C}}$ on $G_{\mathbb{C}}^N$ by diagonal conjugation. For $\underline{a} \in G_{\mathbb{C}}^N$, let $G_{\mathbb{C}} \cdot \underline{a}$ denote the corresponding orbit. Since $G_{\mathbb{C}}$ is not compact, $G_{\mathbb{C}} \cdot \underline{a}$ need not be closed. If a holomorphic function on $G_{\mathbb{C}}^N$ is invariant under the action of G by diagonal conjugation, then it is also invariant under the action of $G_{\mathbb{C}}$ by diagonal conjugation; i.e., it is constant on the orbit $G_{\mathbb{C}} \cdot \underline{a}$ for every $\underline{a} \in G_{\mathbb{C}}^N$. Being continuous, it is then constant on the closure $\overline{G_{\mathbb{C}} \cdot \underline{a}}$. As a consequence, it takes the same value on two orbits whenever their closures intersect. This motivates the following definition. Two elements $\underline{a}, \underline{b} \in G_{\mathbb{C}}^N$ are said to be orbit closure equivalent if there exist $\underline{c}_1, \dots, \underline{c}_r \in G_{\mathbb{C}}^N$ such that

$$\overline{G_{\mathbb{C}} \cdot \underline{a}} \cap \overline{G_{\mathbb{C}} \cdot \underline{c}_1} \neq \emptyset, \quad \overline{G_{\mathbb{C}} \cdot \underline{c}_1} \cap \overline{G_{\mathbb{C}} \cdot \underline{c}_2} \neq \emptyset, \quad \dots, \quad \overline{G_{\mathbb{C}} \cdot \underline{c}_r} \cap \overline{G_{\mathbb{C}} \cdot \underline{b}} \neq \emptyset.$$

Clearly, orbit closure equivalence defines an equivalence relation on $G_{\mathbb{C}}^N$, indeed. Let $G_{\mathbb{C}}^N // G_{\mathbb{C}}$ denote the topological quotient. The notation is motivated by the fact that the quotient provides a categorical quotient of $G_{\mathbb{C}}^N$ by $G_{\mathbb{C}}$ in the sense of geometric invariant theory.³⁰ By construction, the elements of \mathcal{H} descend to continuous functions on $G_{\mathbb{C}}^N // G_{\mathbb{C}}$.

In Ref. 11, we have explained in some detail how the orbit closure quotient $G_{\mathbb{C}}^N // G_{\mathbb{C}}$ is related to the reduced phase space \mathcal{P} . This discussion is based on the observation that, via the equivariant diffeomorphism (3.2), the momentum mapping may be viewed as a mapping

$$\mu: G_{\mathbb{C}}^N \rightarrow \mathfrak{g}^*$$

and, thus, \mathcal{P} may be viewed as the quotient of $\mu^{-1}(0) \subset G_{\mathbb{C}}^N$ by the action of G . In this language, $\mu^{-1}(0)$ turns out to be a Kempf-Ness set.²⁴ Using this fact, one can prove the following.

Theorem 3.2. *The natural inclusion mapping $\mu^{-1}(0) \rightarrow G_{\mathbb{C}}^N$ induces a homeomorphism*

$$\mathcal{P} \rightarrow G_{\mathbb{C}}^N // G_{\mathbb{C}}. \quad (3.4)$$

For the proof, see Ref. 17.

As a by-product of the proof, one finds that two points $\underline{a}, \underline{b} \in G_{\mathbb{C}}^N$ are orbit closure equivalent if and only if

$$\overline{G_{\mathbb{C}} \cdot \underline{a}} \cap \overline{G_{\mathbb{C}} \cdot \underline{b}} \cap \mu^{-1}(0) \neq \emptyset. \quad (3.5)$$

As a result, via the homeomorphism (3.4), the elements of \mathcal{H} can be interpreted as functions on \mathcal{P} . By virtue of this interpretation, to a given orbit type stratum $\mathcal{P}_{\tau} \subset \mathcal{P}$, there corresponds the closed subspace

$$\mathcal{V}_{\tau} := \{\psi \in \mathcal{H} : \psi|_{\mathcal{P}_{\tau}} = 0\}.$$

We define the subspace \mathcal{H}_{τ} associated with \mathcal{P}_{τ} to be the orthogonal complement of \mathcal{V}_{τ} in \mathcal{H} . Then, we have the orthogonal decomposition

$$\mathcal{H}_{\tau} \oplus \mathcal{V}_{\tau} = \mathcal{H}.$$

Remark 3.3. Since holomorphic functions are continuous, one has

$$\mathcal{V}_{\tau} = \{\psi \in \mathcal{H} : \psi|_{\overline{\mathcal{P}_{\tau}}} = 0\}. \quad (3.6)$$

First, since the principal stratum is dense in \mathcal{P} , this implies that the subspace associated with that stratum coincides with \mathcal{H} . Thus, in the discussion of the orbit type subspaces below, the principal stratum may be ignored. Second, recall that in a stratification, the strata satisfy the condition of

the frontier, which means that if $\mathcal{P}_\sigma \cap \overline{\mathcal{P}_\tau} \neq \emptyset$, then $\mathcal{P}_\sigma \subset \overline{\mathcal{P}_\tau}$. In view of this, (3.6) implies that if $\mathcal{P}_\sigma \cap \overline{\mathcal{P}_\tau} \neq \emptyset$, then $\mathcal{V}_\tau \subset \mathcal{V}_\sigma$ and hence $\mathcal{H}_\sigma \subset \mathcal{H}_\tau$. The family of orthogonal projections

$$\mathcal{H}_\tau \rightarrow \mathcal{H}_\sigma \text{ whenever } \mathcal{P}_\sigma \cap \overline{\mathcal{P}_\tau} \neq \emptyset$$

makes the family of closed subspaces \mathcal{H}_τ into a costratification in the sense of Huebschmann.²⁰ ♦

In order to analyze the condition $\psi|_{\mathcal{P}_\tau} = 0$, it is convenient to work with those subsets of $G_{\mathbb{C}}^N$ which under the natural projection $G_{\mathbb{C}}^N \rightarrow G_{\mathbb{C}}^N // G_{\mathbb{C}}$ and the homeomorphism (3.4) correspond to the orbit type strata of \mathcal{P} . For a given orbit type stratum \mathcal{P}_τ , denote this subset by $(G_{\mathbb{C}}^N)_\tau$. That is, $(G_{\mathbb{C}}^N)_\tau$ consists of the elements \underline{a} of $G_{\mathbb{C}}^N$ whose orbit closure equivalence class belongs to the image of \mathcal{P}_τ under the homeomorphism (3.4). In other words, $\underline{a} \in (G_{\mathbb{C}}^N)_\tau$ if and only if it is orbit closure equivalent to some element of $\mu^{-1}(0)$ whose G -orbit belongs to \mathcal{P}_τ . Clearly,

$$\mathcal{V}_\tau = \{\psi \in \mathcal{H} : \psi|_{(G_{\mathbb{C}}^N)_\tau} = 0\}. \quad (3.7)$$

C. Characterization of costrata in terms of relations

To conclude the general discussion, we describe how to construct \mathcal{V}_τ and \mathcal{H}_τ using defining relations for the orbit type strata \mathcal{P}_τ .

Let $\mathfrak{R}(G^N)$ denote the commutative algebra of representative functions on G^N and let $\mathcal{R} := \mathfrak{R}(G^N)^G$ be the subalgebra of G -invariant elements. Since $G_{\mathbb{C}}^N$ is the complexification of the compact Lie group G^N , the proposition and Theorem 3 in Sec. 8.7.2 of Ref. 33 imply that $\mathfrak{R}(G^N)$ coincides with the coordinate ring of $G_{\mathbb{C}}^N$, viewed as a complex affine variety, and that $\mathfrak{R}(G^N)$ coincides with the algebra of representative functions on $G_{\mathbb{C}}^N$. As a consequence, \mathcal{R} coincides with the algebra of G -invariant representative functions on $G_{\mathbb{C}}^N$, where the relation is given by restriction and analytic continuation, respectively.

Recall that an ideal $\mathcal{I} \subset \mathcal{R}$ is called a radical ideal if for all $f \in \mathcal{R}$ satisfying $f^n \in \mathcal{I}$ for some n , one has $f \in \mathcal{I}$. Moreover, given a subset $R \subset \mathcal{R}$, one defines the zero locus of R by

$$\{\underline{a} \in G_{\mathbb{C}}^N : f(\underline{a}) = 0 \text{ for all } f \in R\} \subset G_{\mathbb{C}}^N.$$

It coincides with the zero locus of the ideal in \mathcal{R} generated by R .

Proposition 3.4. Let \mathcal{P}_τ be an orbit type stratum and let R_τ be a subset of \mathcal{R} satisfying

1. The zero locus of R_τ coincides with the topological closure of $(G_{\mathbb{C}}^N)_\tau$.
2. The ideal generated by R_τ in \mathcal{R} is a radical ideal.

Then, \mathcal{V}_τ is obtained by intersecting \mathcal{H} with the ideal generated algebraically by R_τ in the algebra $\text{Hol}(G_{\mathbb{C}}^N)^G$ of G -invariant holomorphic functions on $G_{\mathbb{C}}^N$.

For the proof, see Ref. 11.

By Hilbert's basis theorem, finite subsets $R_\tau \subset \mathcal{R}$ satisfying conditions 1 and 2 of Proposition 3.4 exist. Given R_τ , Proposition 3.4 implies the following explicit characterization of the subspaces \mathcal{V}_τ and \mathcal{H}_τ in terms of multiplication operators. For $f \in \mathcal{R}$, let $\hat{f} : \mathcal{H} \rightarrow \mathcal{H}$ denote the operator of multiplication by f .

Corollary 3.5. Let \mathcal{P}_τ be an orbit type stratum and let $R_\tau = \{p_1, \dots, p_r\}$ be a finite subset of \mathcal{R} satisfying conditions 1 and 2 of Proposition 3.4. Then,

$$\mathcal{V}_\tau = \text{im}(\hat{p}_1) + \dots + \text{im}(\hat{p}_r), \quad \mathcal{H}_\tau = \ker(\hat{p}_1^\dagger) \cap \dots \cap \ker(\hat{p}_r^\dagger). \quad \square$$

In what follows, we will refer to conditions 1 and 2 of Proposition 3.4 as the zero locus condition and the radical ideal condition, respectively.

D. The commutative algebra \mathcal{R}

By Proposition 3.4 and Corollary 3.5, the costratification of the quantum Hilbert space \mathcal{H} is given by a family of finite subsets $R_\tau \subset \mathcal{R}$ satisfying conditions 1 and 2. Each of these subsets consists of a finite set of G -invariant polynomials on $G_{\mathbb{C}}^N$. To construct the costratification explicitly, one has to find the images of the multiplication operators defined by these invariant polynomials. This can be achieved by choosing an orthonormal basis in \mathcal{H} and by finding the structure constants of the multiplication law in \mathcal{R} in that basis.

By Remark 3.1, we can first consider the Hilbert space $L^2(G^N)^G$ and use the theory of compact Lie groups. For the convenience of the reader and to fix the notation, we recall some basics; see, e.g., Ref. 31 or 12 for details. Below, all representations are assumed to be continuous and unitary without further notice. Let \widehat{G} denote the set of isomorphism classes of finite-dimensional irreps of G . Given a finite-dimensional unitary representation (H, π) of G , let $C(G)_\pi \subset \mathfrak{R}(G)$ denote the subspace of representative functions and let $\chi_\pi \in C(G)_\pi$ be the character of π , defined by $\chi_\pi(a) := \text{tr}(\pi(a))$. The same notation will be used for the Lie group G^N .

The elements of \widehat{G} will be labeled by the corresponding highest weight λ relative to some chosen Cartan subalgebra and some chosen dominant Weyl chamber. Assume that for every $\lambda \in \widehat{G}$ a concrete unitary irrep (H_λ, π_λ) of highest weight λ in the Hilbert space H_λ has been chosen. Given $\underline{\lambda} = (\lambda^1, \dots, \lambda^N) \in \widehat{G}^N$, we define a representation $(H_{\underline{\lambda}}, \pi_{\underline{\lambda}})$ of G^N by

$$H_{\underline{\lambda}} = \bigotimes_{i=1}^N H_{\lambda^i}, \quad \pi_{\underline{\lambda}}(a) = \bigotimes_{i=1}^N \pi_{\lambda^i}(a_i), \quad (3.8)$$

where $\underline{a} = (a_1, \dots, a_N)$. This representation is irreducible and we have

$$C(G^N)_{\pi_{\underline{\lambda}}} \cong \bigotimes_{i=1}^N C(G)_{\pi_{\lambda^i}},$$

isometrically with respect to the L^2 -norms. Using this, together with the Peter-Weyl theorem for G , we obtain that $\bigoplus_{\underline{\lambda} \in \widehat{G}^N} C(G^N)_{\pi_{\underline{\lambda}}}$ is dense in $L^2(G^N, d^N a)$. Since $\bigoplus_{\underline{\lambda} \in \widehat{G}^N} C(G^N)_{\pi_{\underline{\lambda}}} \subset \bigoplus_{\pi \in \widehat{G}^N} C(G^N)_\pi$, this implies

Lemma 3.6. Every irreducible representation of G^N is equivalent to a product representation $(H_{\underline{\lambda}}, \pi_{\underline{\lambda}})$ with $\underline{\lambda} \in \widehat{G}^N$. If $(H_{\underline{\lambda}}, \pi_{\underline{\lambda}})$ and $(H_{\underline{\lambda}'}, \pi_{\underline{\lambda}'})$ are isomorphic, then $\underline{\lambda} = \underline{\lambda}'$. \square

Given $\underline{\lambda} \in \widehat{G}^N$, let $\pi_{\underline{\lambda}}^d$ denote the representation of G on $H_{\underline{\lambda}}$ defined by

$$\pi_{\underline{\lambda}}^d(a) := \pi_{\underline{\lambda}}(a, \dots, a). \quad (3.9)$$

This representation will be referred to as the diagonal representation induced by $\pi_{\underline{\lambda}}$. It is reducible and has the isotypical decomposition

$$H_{\underline{\lambda}} = \bigoplus_{\lambda \in \widehat{G}} H_{\underline{\lambda}, \lambda}$$

into uniquely determined subspaces $H_{\underline{\lambda}, \lambda}$. Recall that these subspaces may be obtained as the images of the orthogonal projectors

$$\mathbb{P}_\lambda := \dim(H_\lambda) \int_G \overline{\chi_{\pi_\lambda}(a)} \pi_{\underline{\lambda}}(a) da \quad (3.10)$$

on $H_{\underline{\lambda}}$. These projectors commute with one another and with $\pi_{\underline{\lambda}}^d$. If an isotypical subspace $H_{\underline{\lambda}, \lambda}$ is reducible, we can further decompose it in a non-unique way into irreducible subspaces of isomorphism type λ . Let $m_{\underline{\lambda}}(\lambda)$ denote the number of these irreducible subspaces (the multiplicity of π_λ in $\pi_{\underline{\lambda}}^d$)

and let $\widehat{G}_{\underline{\lambda}}$ denote the subset of \widehat{G} consisting of the highest weights λ such that $m_{\underline{\lambda}}(\lambda) > 0$. This way, we obtain a unitary G -representation isomorphism

$$\varphi_{\underline{\lambda}} : H_{\underline{\lambda}} \rightarrow \bigoplus_{\lambda \in \widehat{G}_{\underline{\lambda}}} \bigoplus_{k=1}^{m_{\underline{\lambda}}(\lambda)} H_{\lambda}. \quad (3.11)$$

Let

$$\text{pr}_{\underline{\lambda},k}^{\underline{\lambda}} : \bigoplus_{\lambda \in \widehat{G}_{\underline{\lambda}}} \bigoplus_{k=1}^{m_{\underline{\lambda}}(\lambda)} H_{\lambda} \rightarrow H_{\underline{\lambda}}, \quad \text{i}_{\underline{\lambda},k}^{\underline{\lambda}} : H_{\underline{\lambda}} \rightarrow \bigoplus_{\lambda \in \widehat{G}_{\underline{\lambda}}} \bigoplus_{k=1}^{m_{\underline{\lambda}}(\lambda)} H_{\lambda},$$

denote the natural projections and injections of the direct sum, respectively. For every $\lambda \in \widehat{G}_{\underline{\lambda}}$ and every $k, l = 1, \dots, m_{\underline{\lambda}}(\lambda)$, define a G -representation endomorphism $A_{k,l}^{\underline{\lambda},\lambda}$ of $\pi_{\underline{\lambda}}^d$ by

$$A_{k,l}^{\underline{\lambda},\lambda} := \frac{1}{\sqrt{\dim(H_{\underline{\lambda}})}} \varphi_{\underline{\lambda}}^{-1} \circ \text{i}_{\underline{\lambda},k}^{\underline{\lambda}} \circ \text{pr}_{\underline{\lambda},l}^{\underline{\lambda}} \circ \varphi_{\underline{\lambda}} \quad (3.12)$$

and a G -invariant function $(\chi_{\underline{\lambda}})_{k,l}^{\lambda}$ on G^N by

$$(\chi_{\underline{\lambda}})_{k,l}^{\lambda}(\underline{a}) := \sqrt{\dim(H_{\underline{\lambda}})} \text{tr}(\pi_{\underline{\lambda}}(\underline{a}) A_{k,l}^{\underline{\lambda},\lambda}). \quad (3.13)$$

Proposition 3.7. The family of functions

$$\{(\chi_{\underline{\lambda}})_{k,l}^{\lambda} : \underline{\lambda} \in \widehat{G}^N, \lambda \in \widehat{G}_{\underline{\lambda}}, k, l = 1, \dots, m_{\underline{\lambda}}(\lambda)\}$$

constitutes an orthonormal basis in $L^2(G^N)^G$.

Proof. Note that for every $\underline{\lambda} \in \widehat{G}^N$, the mapping

$$T_{\underline{\lambda}} : \text{End}(H_{\underline{\lambda}}) \rightarrow C(G^N)_{\pi_{\underline{\lambda}}}, \quad T_{\underline{\lambda}}(A)(a) := \sqrt{\dim H_{\underline{\lambda}}} \text{tr}(\pi_{\underline{\lambda}}(a)A), \quad (3.14)$$

is a unitary G -representation isomorphism with respect to the scalar product on $\text{End}(H_{\underline{\lambda}})$ defined by $\langle A|B \rangle = \text{tr}(A^*B)$ and the induced endomorphism representation on $\text{End}(H_{\underline{\lambda}})$, given by assigning to $g \in G$ the automorphism

$$A \mapsto \pi_{\underline{\lambda}}^d(g) A \pi_{\underline{\lambda}}^d(g)^{-1}$$

of $\text{End}(H)$. Being a representation isomorphism, $T_{\underline{\lambda}}$ restricts to a unitary Hilbert space isomorphism of the subspaces of G -invariant elements, $\text{End}(H_{\underline{\lambda}})^G \rightarrow C(G^N)_{\pi_{\underline{\lambda}}}^G$. Now, $\text{End}(H_{\underline{\lambda}})^G$ consists precisely of the representation endomorphisms of $\pi_{\underline{\lambda}}^d$. Hence, Schur's lemma implies that it is spanned by the endomorphisms $A_{k,l}^{\underline{\lambda},\lambda}$ with $\lambda \in \widehat{G}_{\underline{\lambda}}$ and $k, l = 1, \dots, m_{\underline{\lambda}}(\lambda)$. Using $\text{pr}_{\underline{\lambda},k}^{\underline{\lambda}} \circ \text{i}_{\underline{\lambda}',k'}^{\underline{\lambda}} = \delta_{\lambda\lambda'} \delta_{kk'} \text{id}_{H_{\underline{\lambda}}}$, we compute

$$\langle A_{k,l}^{\underline{\lambda},\lambda} | A_{k',l'}^{\underline{\lambda},\lambda'} \rangle = \delta_{\lambda\lambda'} \delta_{kk'} \delta_{ll'}.$$

It follows that the endomorphisms $A_{k,l}^{\underline{\lambda},\lambda}$ with $\lambda \in \widehat{G}_{\underline{\lambda}}$ and $k, l = 1, \dots, m_{\underline{\lambda}}(\lambda)$ form an orthonormal basis in $\text{End}(H_{\underline{\lambda}})^G$ and hence that their images under $T_{\underline{\lambda}}$, i.e., the functions $(\chi_{\underline{\lambda}})_{k,l}^{\lambda}$, form an orthonormal basis in $C(G^N)_{\pi_{\underline{\lambda}}}^G$. Thus, the family given in the proposition yields an orthonormal basis in $\mathcal{R} = \mathfrak{R}(G^N)^G$.

It remains to show that \mathcal{R} is dense in $L^2(G^N, \text{d}^N a)^G$. This follows from the Peter-Weyl theorem for G^N by applying the averaging operator

$$\mathbb{P}_G : L^2(G^N, d^N a) \rightarrow L^2(G^N, d^N a)^G, \quad \mathbb{P}_G(f)(\underline{a}) = \int_G f(ga_1g^{-1}, \dots, ga_Ng^{-1}) dg$$

and observing that the image of a dense subset under a surjective continuous mapping is dense. \square

By analytic continuation, the irreps π_λ of G induce irreps $\pi_\lambda^\mathbb{C}$ of $G_\mathbb{C}$, the irreps $\pi_{\underline{\lambda}}$ of G^N induce irreps $\pi_{\underline{\lambda}}^\mathbb{C}$ of $G_\mathbb{C}^N$, and the functions $(\chi_{\underline{\lambda}})_k^\lambda$ on G^N induce holomorphic functions $(\chi_{\underline{\lambda}}^\mathbb{C})_{k,l}^\lambda$ on $G_\mathbb{C}^N$. Then, (3.8), (3.9), and (3.13) hold with $\pi_{\underline{\lambda}}$, π_λ , and $(\chi_{\underline{\lambda}})_k^\lambda$ replaced by, respectively, $\pi_{\underline{\lambda}}^\mathbb{C}$, $\pi_\lambda^\mathbb{C}$, and $(\chi_{\underline{\lambda}}^\mathbb{C})_{k,l}^\lambda$.

Corollary 3.8. The family of functions

$$\{(\chi_{\underline{\lambda}}^\mathbb{C})_{k,l}^\lambda : \underline{\lambda} \in \widehat{G}^N, \lambda \in \widehat{G}_{\underline{\lambda}}, k, l = 1, \dots, m_{\underline{\lambda}}(\lambda)\}$$

constitutes an orthogonal basis in \mathcal{H} . The norms are

$$\|(\chi_{\underline{\lambda}}^\mathbb{C})_{k,l}^\lambda\|^2 = \prod_{r=1}^N C_{\lambda^r}, \quad C_{\lambda^r} = (\hbar\pi)^{\dim(G)/2} e^{\hbar|\lambda^r + \rho|^2}, \quad (3.15)$$

where ρ denotes half the sum of the positive roots. The expansion coefficients of $f \in \mathcal{H}$ with respect to this basis are given by the scalar products $\langle (\chi_{\underline{\lambda}}^\mathbb{C})_{k,l}^\lambda | f|_{G^N} \rangle$ in $L^2(G^N)^G$.

Proof. See [Appendix A](#). The last statement follows from the fact that two elements of \mathcal{H} coincide if and only if their restrictions to G^N coincide. Since the functions $(\chi_{\underline{\lambda}})_k^\lambda$ form an orthonormal basis in $L^2(G^N)^G$, we have $f|_{G^N} = \sum_{\underline{\lambda}, \lambda, k, l} \langle (\chi_{\underline{\lambda}})_k^\lambda | f|_{G^N} \rangle (\chi_{\underline{\lambda}})_k^\lambda$. Since $(\chi_{\underline{\lambda}})_k^\lambda = (\chi_{\underline{\lambda}}^\mathbb{C})_{k,l}^\lambda|_{G^N}$, this yields the assertion. \square

Remark 3.9. The orthonormal basis of invariant representative functions provided by Proposition 3.7 is a special case of a spin network basis in the sense of Baez.³ It is special in so far as from the very beginning we have fixed a tree gauge, which reduces the group of local gauge transformations to the action of G . Moreover, our basis above clearly corresponds to a fixed graph (a finite regular cubic lattice). In this situation, we are able to provide a more explicit presentation of the basis elements in terms of appropriate functions. We refer to Ref. 3 for comments on various applications of spin networks in mathematical physics. In particular, over the years, spin network states have become an important tool in loop quantum gravity; see Ref. 39 and further references therein. \blacklozenge

Now, let us turn to the discussion of the multiplication structure of the G -invariant representative functions $(\chi_{\underline{\lambda}}^\mathbb{C})_{k,l}^\lambda$. We assume that a unitary G -representation isomorphism (3.11) has been chosen for every $\underline{\lambda} \in \widehat{G}^N$ and every N . Denote

$$d_\lambda := \dim H_\lambda, \quad d_{\underline{\lambda}} := \dim H_{\underline{\lambda}}.$$

Writing

$$\begin{aligned} & (\chi_{\underline{\lambda}_1}^\mathbb{C})_{k_1, l_1}^{\lambda_1}(\underline{a}) (\chi_{\underline{\lambda}_2}^\mathbb{C})_{k_2, l_2}^{\lambda_2}(\underline{a}) \\ &= \sqrt{d_{\underline{\lambda}_1} d_{\underline{\lambda}_2}} \operatorname{tr} \left(\left(A_{k_1, l_1}^{\lambda_1, \lambda_1} \otimes A_{k_2, l_2}^{\lambda_2, \lambda_2} \right) \circ \left(\pi_{\underline{\lambda}_1}(\underline{a}) \otimes \pi_{\underline{\lambda}_2}(\underline{a}) \right) \right), \end{aligned} \quad (3.16)$$

we see that in order to expand the product $(\chi_{\underline{\lambda}_1}^\mathbb{C})_{k_1, l_1}^{\lambda_1} \cdot (\chi_{\underline{\lambda}_2}^\mathbb{C})_{k_2, l_2}^{\lambda_2}$ in terms of the basis functions $(\chi_{\underline{\lambda}}^\mathbb{C})_{k,l}^\lambda$, a reasonable strategy is to decompose the G^N -representation $\pi_{\underline{\lambda}_1} \otimes \pi_{\underline{\lambda}_2}$ into G^N -irreps $\underline{\lambda}$ and then relate these G^N -irreps to the basis functions using the chosen G -representation isomorphisms $\varphi_{\underline{\lambda}}$. To implement this, we define two different unitary G -representation isomorphisms of the diagonal representation $\pi_{\underline{\lambda}_1}^d \otimes \pi_{\underline{\lambda}_2}^d$ with an orthogonal direct sum of G -irreps. The first one, $\Phi_{\underline{\lambda}_1, \underline{\lambda}_2}$, is adapted to the tensor product on the right-hand side of (3.16). It is defined by

$$\Phi_{\underline{\lambda}_1, \underline{\lambda}_2} : H_{\underline{\lambda}_1} \otimes H_{\underline{\lambda}_2} \xrightarrow{\varphi_{\underline{\lambda}_1} \otimes \varphi_{\underline{\lambda}_2}} \bigoplus_{\underline{\lambda}_1, \underline{\lambda}_2} \left(\bigoplus_{i_1=1}^{m_{\underline{\lambda}_1}(\underline{\lambda}_1)} \bigoplus_{i_2=1}^{m_{\underline{\lambda}_2}(\underline{\lambda}_2)} (H_{\underline{\lambda}_1} \otimes H_{\underline{\lambda}_2}) \right) \quad (3.17)$$

$$\xrightarrow{\phi_{\underline{\lambda}_1, \underline{\lambda}_2}} \bigoplus_{\underline{\lambda}_1, \underline{\lambda}_2} \left(\bigoplus_{i_1=1}^{m_{\underline{\lambda}_1}(\underline{\lambda}_1)} \bigoplus_{i_2=1}^{m_{\underline{\lambda}_2}(\underline{\lambda}_2)} \left(\bigoplus_{\underline{\lambda}} \bigoplus_{i=1}^{m_{\underline{\lambda}_1, \underline{\lambda}_2}(\underline{\lambda})} H_{\underline{\lambda}} \right) \right), \quad (3.18)$$

where $\phi_{\underline{\lambda}_1, \underline{\lambda}_2}$ acts on each summand $H_{\underline{\lambda}_1} \otimes H_{\underline{\lambda}_2}$ as $\varphi_{(\underline{\lambda}_1, \underline{\lambda}_2)}$. Let $\text{pr}_{\underline{\lambda}_1, \underline{\lambda}_2, i_1, i_2, \underline{\lambda}, i}^{\underline{\lambda}_1, \underline{\lambda}_2}$ and $i_{\underline{\lambda}_1, \underline{\lambda}_2, i_1, i_2, \underline{\lambda}, i}^{\underline{\lambda}_1, \underline{\lambda}_2}$ denote the natural projection and injection operators of the direct sum (3.18), respectively. We have

$$\varphi_{(\underline{\lambda}_1, \underline{\lambda}_2)} \circ \left(\text{pr}_{\underline{\lambda}_1, i_1}^{\underline{\lambda}_1} \otimes \text{pr}_{\underline{\lambda}_2, i_2}^{\underline{\lambda}_2} \right) = \left(\sum_{\underline{\lambda}} \sum_{i=1}^{m_{(\underline{\lambda}_1, \underline{\lambda}_2)}(\underline{\lambda})} i_{\underline{\lambda}, i}^{(\underline{\lambda}_1, \underline{\lambda}_2)} \circ \text{pr}_{\underline{\lambda}_1, \underline{\lambda}_2, i_1, i_2, \underline{\lambda}, i}^{\underline{\lambda}_1, \underline{\lambda}_2} \right) \circ \phi_{\underline{\lambda}_1, \underline{\lambda}_2}, \quad (3.19)$$

$$\phi_{\underline{\lambda}_1, \underline{\lambda}_2} \circ \left(i_{\underline{\lambda}_1, i_1}^{\underline{\lambda}_1} \otimes i_{\underline{\lambda}_2, i_2}^{\underline{\lambda}_2} \right) = \left(\sum_{\underline{\lambda}} \sum_{i=1}^{m_{(\underline{\lambda}_1, \underline{\lambda}_2)}(\underline{\lambda})} i_{\underline{\lambda}_1, \underline{\lambda}_2, i_1, i_2, \underline{\lambda}, i}^{\underline{\lambda}_1, \underline{\lambda}_2} \circ \text{pr}_{\underline{\lambda}, i}^{(\underline{\lambda}_1, \underline{\lambda}_2)} \right) \circ \varphi_{(\underline{\lambda}_1, \underline{\lambda}_2)}. \quad (3.20)$$

The second unitary G -representation isomorphism, $\Psi_{\underline{\lambda}_1, \underline{\lambda}_2}$, is adapted to the definition of the basis functions $(\chi_{\underline{\lambda}})_{k, l}^l$. It is defined by

$$\Psi_{\underline{\lambda}_1, \underline{\lambda}_2} : H_{\underline{\lambda}_1} \otimes H_{\underline{\lambda}_2} \xrightarrow{\psi_{\underline{\lambda}_1, \underline{\lambda}_2}^1} \bigoplus_{\underline{\lambda}} \left(\bigoplus_{i=1}^{m_{\underline{\lambda}_1, \underline{\lambda}_2}(\underline{\lambda})} H_{\underline{\lambda}} \right) \quad (3.21)$$

$$\xrightarrow{\psi_{\underline{\lambda}_1, \underline{\lambda}_2}^2} \bigoplus_{\underline{\lambda}} \left(\bigoplus_{i=1}^{m_{\underline{\lambda}_1, \underline{\lambda}_2}(\underline{\lambda})} \left(\bigoplus_{\underline{\lambda}} \bigoplus_{k=1}^{m_{\underline{\lambda}}(\underline{\lambda})} H_{\underline{\lambda}} \right) \right), \quad (3.22)$$

where $\psi_{\underline{\lambda}_1, \underline{\lambda}_2}^1$ is some unitary G^N -representation isomorphism, provided by Lemma 3.6, and $\psi_{\underline{\lambda}_1, \underline{\lambda}_2}^2$ is the G -representation isomorphism acting on each summand $H_{\underline{\lambda}}$ as $\varphi_{\underline{\lambda}}$. Moreover, $m_{\underline{\lambda}_1, \underline{\lambda}_2}(\underline{\lambda})$ is the multiplicity of the G^N -irrep $H_{\underline{\lambda}}$ in $H_{\underline{\lambda}_1} \otimes H_{\underline{\lambda}_2}$. Let $\text{pr}_{\underline{\lambda}, i}^{\underline{\lambda}_1, \underline{\lambda}_2}$ and $i_{\underline{\lambda}, i}^{\underline{\lambda}_1, \underline{\lambda}_2}$ be the natural projections and injections, respectively, of the direct sum (3.21), and let $\text{pr}_{\underline{\lambda}, i, \underline{\lambda}, k}^{\underline{\lambda}_1, \underline{\lambda}_2}$ and $i_{\underline{\lambda}, i, \underline{\lambda}, k}^{\underline{\lambda}_1, \underline{\lambda}_2}$ be the natural projections and injections, respectively, of the direct sum (3.22), respectively. We have

$$\psi_{\underline{\lambda}_1, \underline{\lambda}_2}^2 \circ i_{\underline{\lambda}, i}^{\underline{\lambda}_1, \underline{\lambda}_2} = \left(\sum_{\underline{\lambda}} \sum_{k=1}^{m_{\underline{\lambda}}(\underline{\lambda})} i_{\underline{\lambda}, i, \underline{\lambda}, k}^{\underline{\lambda}_1, \underline{\lambda}_2} \circ \text{pr}_{\underline{\lambda}, k}^{\underline{\lambda}} \right) \circ \varphi_{\underline{\lambda}}, \quad (3.23)$$

$$\varphi_{\underline{\lambda}} \circ \text{pr}_{\underline{\lambda}, i}^{\underline{\lambda}_1, \underline{\lambda}_2} = \left(\sum_{\underline{\lambda}} \sum_{k=1}^{m_{\underline{\lambda}}(\underline{\lambda})} i_{\underline{\lambda}, k}^{\underline{\lambda}} \circ \text{pr}_{\underline{\lambda}, i, \underline{\lambda}, k}^{\underline{\lambda}_1, \underline{\lambda}_2} \right) \circ \psi_{\underline{\lambda}_1, \underline{\lambda}_2}^2. \quad (3.24)$$

By construction, $\Psi_{\underline{\lambda}_1, \underline{\lambda}_2} \circ \Phi_{\underline{\lambda}_1, \underline{\lambda}_2}^{-1}$ is a unitary automorphism of a direct sum of G -irreps $H_{\underline{\lambda}}$. Hence, Schur's lemma implies that

$$\text{pr}_{\underline{\lambda}, i, \underline{\lambda}, k}^{\underline{\lambda}_1, \underline{\lambda}_2} \circ \left(\Psi_{\underline{\lambda}_1, \underline{\lambda}_2} \circ \Phi_{\underline{\lambda}_1, \underline{\lambda}_2}^{-1} \right) \circ i_{\underline{\lambda}_1, \underline{\lambda}_2, i_1, i_2, \underline{\lambda}', i'}^{\underline{\lambda}_1, \underline{\lambda}_2} = \delta_{\underline{\lambda} \underline{\lambda}'} U_{\underline{\lambda}_1, \underline{\lambda}_2, i_1, i_2, i', i'}^{\underline{\lambda}_1, \underline{\lambda}_2, i, i, \underline{\lambda}, k} \text{id}_{H_{\underline{\lambda}}}, \quad (3.25)$$

$$\text{pr}_{\underline{\lambda}_1, \underline{\lambda}_2, i_1, i_2, \underline{\lambda}', i'}^{\underline{\lambda}_1, \underline{\lambda}_2} \circ \left(\Phi_{\underline{\lambda}_1, \underline{\lambda}_2} \circ \Psi_{\underline{\lambda}_1, \underline{\lambda}_2}^{-1} \right) \circ i_{\underline{\lambda}, i, \underline{\lambda}, k}^{\underline{\lambda}_1, \underline{\lambda}_2} = \delta_{\underline{\lambda} \underline{\lambda}'} \overline{U_{\underline{\lambda}_1, \underline{\lambda}_2, i_1, i_2, i', i'}^{\underline{\lambda}_1, \underline{\lambda}_2, i, i, \underline{\lambda}, k}} \text{id}_{H_{\underline{\lambda}}}, \quad (3.26)$$

with certain coefficients $U_{\underline{\lambda}_1, \underline{\lambda}_2, i_1, i_2, i', i'}^{\underline{\lambda}_1, \underline{\lambda}_2, i, i, \underline{\lambda}, k}$.

Proposition 3.10. In terms of the basis functions, the multiplication in \mathcal{R} is given by

$$\begin{aligned}
& (\chi_{\underline{\lambda}_1}^{\mathbb{C}})^{\lambda_1}_{k_1, l_1} \cdot (\chi_{\underline{\lambda}_2}^{\mathbb{C}})^{\lambda_2}_{k_2, l_2} \\
&= \sqrt{\frac{d_{\underline{\lambda}_1} d_{\underline{\lambda}_2}}{d_{\lambda_1} d_{\lambda_2}}} \sum_{\underline{\lambda}} \sum_{n=1}^{m_{\underline{\lambda}_1, \underline{\lambda}_2}(\underline{\lambda})} \sum_{\lambda} \sum_{k, l=1}^{m_{\lambda}(\lambda)} \sum_{j=1}^{m_{(\lambda_1, \lambda_2)}(\lambda)} \sqrt{d_{\underline{\lambda}}} U_{\lambda_1, \lambda_2, k_1, k_2, j}^{\lambda_1, \lambda_2, \underline{\lambda}, n, \lambda, k} \overline{U_{\lambda_1, \lambda_2, l_1, l_2, j}^{\lambda_1, \lambda_2, \underline{\lambda}, n, \lambda, l}} (\chi_{\underline{\lambda}}^{\mathbb{C}})^{\lambda}_{k, l}.
\end{aligned}$$

The same formula holds true for the basis functions $(\chi_{\underline{\lambda}})^{\lambda}_{k, l}$ on G^N .

Proof. It suffices to prove the assertion for the basis functions $(\chi_{\underline{\lambda}})^{\lambda}_{k, l}$ on G^N . In the proof, we will use the shorthand notation

$$z \equiv \frac{1}{\sqrt{d_{\underline{\lambda}_1} d_{\underline{\lambda}_2}}} (\chi_{\underline{\lambda}_1})^{\lambda_1}_{k_1, l_1}(\underline{a}) (\chi_{\underline{\lambda}_2})^{\lambda_2}_{k_2, l_2}(\underline{a}), \quad \Phi \equiv \Phi_{\underline{\lambda}_1, \underline{\lambda}_2}, \quad \Psi \equiv \Psi_{\underline{\lambda}_1, \underline{\lambda}_2}.$$

Using (3.16) and the fact that $\psi_{\underline{\lambda}_1, \underline{\lambda}_2}^1$ is a G^N -representation isomorphism, we may rewrite

$$z = \text{tr} \left(\left(\bigoplus_{\underline{\lambda}} \bigoplus_{i=1}^{m_{\underline{\lambda}_1, \underline{\lambda}_2}(\underline{\lambda})} \pi_{\underline{\lambda}}(\underline{a}) \right) \circ \psi_{\underline{\lambda}_1, \underline{\lambda}_2}^1 \circ \left(A_{k_1, l_1}^{\underline{\lambda}_1, \lambda_1} \otimes A_{k_2, l_2}^{\underline{\lambda}_2, \lambda_2} \right) \circ (\psi_{\underline{\lambda}_1, \underline{\lambda}_2}^1)^{-1} \right).$$

Since

$$\bigoplus_{\underline{\lambda}} \bigoplus_{i=1}^{m_{\underline{\lambda}_1, \underline{\lambda}_2}(\underline{\lambda})} \pi_{\underline{\lambda}}(\underline{a}) = \sum_{\underline{\lambda}} \sum_{i=1}^{m_{\underline{\lambda}_1, \underline{\lambda}_2}(\underline{\lambda})} i_{\underline{\lambda}, i}^{\underline{\lambda}_1, \underline{\lambda}_2} \circ \pi_{\underline{\lambda}}(\underline{a}) \circ \text{pr}_{\underline{\lambda}, i}^{\underline{\lambda}_1, \underline{\lambda}_2},$$

this can be further rewritten as

$$z = \sum_{\underline{\lambda}} \sum_{i=1}^{m_{\underline{\lambda}_1, \underline{\lambda}_2}(\underline{\lambda})} \text{tr} \left(\pi_{\underline{\lambda}}(\underline{a}) \circ \text{pr}_{\underline{\lambda}, i}^{\underline{\lambda}_1, \underline{\lambda}_2} \circ \psi_{\underline{\lambda}_1, \underline{\lambda}_2}^1 \circ \left(A_{k_1, l_1}^{\underline{\lambda}_1, \lambda_1} \otimes A_{k_2, l_2}^{\underline{\lambda}_2, \lambda_2} \right) \circ (\psi_{\underline{\lambda}_1, \underline{\lambda}_2}^1)^{-1} \circ i_{\underline{\lambda}, i}^{\underline{\lambda}_1, \underline{\lambda}_2} \right).$$

By (3.23) and (3.24),

$$\begin{aligned}
i_{\underline{\lambda}, i}^{\underline{\lambda}_1, \underline{\lambda}_2} &= \sum_{\lambda} \sum_{k=1}^{m_{\lambda}(\lambda)} (\psi_{\underline{\lambda}_1, \underline{\lambda}_2}^2)^{-1} \circ i_{\underline{\lambda}, i, \lambda, k}^{\underline{\lambda}_1, \underline{\lambda}_2} \circ \text{pr}_{\lambda, k}^{\underline{\lambda}} \circ \varphi_{\underline{\lambda}}, \\
\text{pr}_{\underline{\lambda}, i}^{\underline{\lambda}_1, \underline{\lambda}_2} &= \sum_{\lambda} \sum_{k=1}^{m_{\lambda}(\lambda)} \varphi_{\underline{\lambda}}^{-1} \circ i_{\lambda, k}^{\underline{\lambda}} \circ \text{pr}_{\underline{\lambda}, i, \lambda, k}^{\underline{\lambda}_1, \underline{\lambda}_2} \circ \psi_{\underline{\lambda}_1, \underline{\lambda}_2}^2.
\end{aligned}$$

Plugging this in, we obtain

$$\begin{aligned}
z &= \sum_{\underline{\lambda}} \sum_{i=1}^{m_{\underline{\lambda}_1, \underline{\lambda}_2}(\underline{\lambda})} \sum_{\lambda, \lambda'} \sum_{k=1}^{m_{\lambda}(\lambda)} \sum_{l=1}^{m_{\lambda}(\lambda')} \text{tr} \left(\pi_{\underline{\lambda}}(\underline{a}) \circ \varphi_{\underline{\lambda}}^{-1} \circ i_{\lambda, k}^{\underline{\lambda}} \circ \text{pr}_{\underline{\lambda}, i, \lambda, k}^{\underline{\lambda}_1, \underline{\lambda}_2} \circ (\Psi \circ \Phi^{-1}) \right. \\
&\quad \left. \circ \Phi \circ \left(A_{k_1, l_1}^{\underline{\lambda}_1, \lambda_1} \otimes A_{k_2, l_2}^{\underline{\lambda}_2, \lambda_2} \right) \circ \Phi^{-1} \circ (\Phi \circ \Psi^{-1}) \circ i_{\underline{\lambda}, i, \lambda, k}^{\underline{\lambda}_1, \underline{\lambda}_2} \circ \text{pr}_{\lambda, k}^{\underline{\lambda}} \circ \varphi_{\underline{\lambda}} \right).
\end{aligned}$$

Using (3.19) and (3.20), we find

$$\Phi \circ \left(A_{k_1, l_1}^{\underline{\lambda}_1, \lambda_1} \otimes A_{k_2, l_2}^{\underline{\lambda}_2, \lambda_2} \right) \circ \Phi^{-1} = \frac{1}{\sqrt{d_{\lambda_1} d_{\lambda_2}}} \sum_{\lambda''} \sum_{j=1}^{m_{(\lambda_1, \lambda_2)}(\lambda'')} i_{\lambda_1, \lambda_2, k_1, k_2, \lambda'' j}^{\underline{\lambda}_1, \underline{\lambda}_2} \circ \text{pr}_{\lambda_1, \lambda_2, l_1, l_2, \lambda'' j}^{\underline{\lambda}_1, \underline{\lambda}_2}.$$

Together with (3.25) and (3.26), this yields, after taking the sums over λ' and λ'' ,

$$\begin{aligned}
z &= \frac{1}{\sqrt{d_{\lambda_1} d_{\lambda_2}}} \sum_{\underline{\lambda}} \sum_{i=1}^{m_{\underline{\lambda}_1, \underline{\lambda}_2}(\underline{\lambda})} \sum_{\lambda} \sum_{k, l=1}^{m_{\lambda}(\lambda)} \sum_{j=1}^{m_{(\lambda_1, \lambda_2)}(\lambda)} U_{\lambda_1, \lambda_2, k_1, k_2, j}^{\underline{\lambda}_1, \underline{\lambda}_2, \underline{\lambda}, i, \lambda, k} \overline{U_{\lambda_1, \lambda_2, l_1, l_2, j}^{\underline{\lambda}_1, \underline{\lambda}_2, \underline{\lambda}, i, \lambda, l}} \\
&\quad \text{tr} \left(\pi_{\underline{\lambda}}(\underline{a}) \circ \varphi_{\underline{\lambda}}^{-1} \circ i_{\lambda, k}^{\underline{\lambda}} \circ \text{pr}_{\lambda, k}^{\underline{\lambda}} \circ \varphi_{\underline{\lambda}} \right).
\end{aligned}$$

The assertion now follows from (3.12) and (3.13). \square

Remark 3.11. Note that the coefficients U in Proposition 3.10 depend on the unitary G -representation isomorphisms $\Phi_{\underline{\lambda}_1, \underline{\lambda}_2}$ and $\Psi_{\underline{\lambda}_1, \underline{\lambda}_2}$. In Subsection IV A, we will see that for $G = \text{SU}(2)$, these isomorphisms are uniquely determined by the choice of a unitary G -representation isomorphism $\varphi_{\underline{\lambda}}$ for every $\underline{\lambda} \in \widehat{G}^N$ and that the coefficients U boil down to recoupling coefficients of angular momentum theory. \blacklozenge

IV. THE MODEL FOR $G = \text{SU}(2)$

A. The commutative algebra \mathcal{R} for $G = \text{SU}(2)$

As observed in Sec. III, to fix concrete basis functions $(\chi_{\underline{\lambda}})_{ij}^{\underline{\lambda}}$, we have to fix the unitary G -representation isomorphisms $\varphi_{\underline{\lambda}}$ entering their definition. As a consequence, we obtain concrete formulae for the unitary operators in the multiplication law of the above algebra, expressed in terms of $\text{SU}(2)$ -recoupling coefficients. This relates the algebra structure to the combinatorics of recoupling theory of angular momentum as provided in Refs. 4, 5, 28, and 41.

In the case of $G = \text{SU}(2)$, the highest weights λ of irreps correspond 1–1 to spins $j = 0, \frac{1}{2}, 1, \frac{3}{2}, \dots$. We will use the common notation D^j for π_j . Thus, (H_j, D^j) is the standard $\text{SU}(2)$ -irrep of spin j , spanned by the orthonormal ladder basis $\{|j, m\rangle : m = -j, -j+1, \dots, j\}$ which is unique up to a phase. Accordingly, every sequence $\underline{\lambda}$ of highest weights corresponds to a sequence \underline{j} of spins. We write $D^{\underline{j}} \equiv \pi_{\underline{j}}$ for the corresponding irrep of $\text{SU}(2)^N$ and $D_{\underline{d}}^{\underline{j}} \equiv \pi_{\underline{d}}^{\underline{j}}$ for the induced diagonal representation of $\text{SU}(2)$. To fix the G -representation isomorphisms

$$\varphi_{\underline{j}} : H_{\underline{j}} \rightarrow \bigoplus_j \bigoplus_{i=1}^{m_j(j)} H_j, \quad (4.1)$$

we choose the following reduction scheme for tensor products of N irreps of $\text{SU}(2)$. Given nonnegative half integers s_1, s_2 , denote

$$\langle s_1, s_2 \rangle := \{|s_1 - s_2|, |s_1 - s_2| + 1, |s_1 - s_2| + 2, \dots, s_1 + s_2\}$$

and recall that the representation space $H_{s_1} \otimes H_{s_2}$ decomposes into unique irreducible subspaces $(H_{s_1} \otimes H_{s_2})_s$ of spin $s \in \langle s_1, s_2 \rangle$. We start with decomposing $H_{j^1} \otimes H_{j^2}$ into the unique irreducible subspaces $(H_{j^1} \otimes H_{j^2})_{l^2}$ with $l^2 \in \langle j^1, j^2 \rangle$. Then, we decompose the invariant subspaces

$$(H_{j^1} \otimes H_{j^2})_{l^2} \otimes H_{j^3} \subset H_{j^1} \otimes H_{j^2} \otimes H_{j^3}$$

into unique irreducible subspaces

$$((H_{j^1} \otimes H_{j^2})_{l^2} \otimes H_{j^3})_{l^3}, \quad l^3 \in \langle l^2, j^3 \rangle.$$

Iterating this, we end up with a decomposition of $H_{\underline{j}}$ into unique irreducible subspaces

$$H_{\underline{j}, \underline{l}} := (\cdots ((H_{j^1} \otimes H_{j^2})_{l^2} \otimes H_{j^3})_{l^3} \cdots \otimes H_{j^N})_{l^N}, \quad (4.2)$$

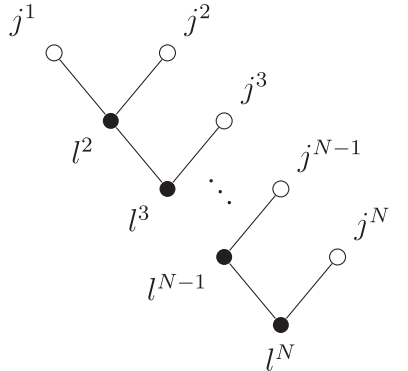
where $\underline{l} = (l^1, \dots, l^N)$ is a sequence of nonnegative half integers satisfying $l^1 = j^1$ and $l^i \in \langle l^{i-1}, j^i \rangle$ for $i = 2, 3, \dots, N$. Let us denote the totality of such sequences by $R(\underline{j})$. Moreover, denote

$$\langle \underline{j} \rangle := \{ \underline{j} : \exists \underline{l} \in R(\underline{j}) \text{ s. th. } j = l^N \}, \quad R(\underline{j}, j) = \{ \underline{l} \in R(\underline{j}) : l^N = j \}.$$

Then, $m_{\underline{j}}(j) = |R(\underline{j}, j)|$ and hence $m_{\underline{j}}(j) \neq 0$ if and only if $j \in \langle \underline{j} \rangle$, and the isotypical component of type j of $H_{\underline{j}}$ is given by the direct sum of the subspaces $H_{\underline{j}, \underline{l}}$ with $\underline{l} \in R(\underline{j}, j)$.

Remark 4.1. Reduction schemes for N -fold tensor products of $\text{SU}(2)$ -irreps of spins j^1, \dots, j^N can be visualized by binary trees with N terminal points \circ labeled by j^1, \dots, j^N and with $N - 1$ internal points \bullet which have two incoming lines and, except for the last one, one outgoing line. The terminal points represent the tensor factors and the internal points represent the intermediate reduction steps

given by the irreducible subspaces in the tensor product of the incoming irreps. The last internal point represents the final irreducible subspace obtained by the reduction scheme. Every labeling of the internal points which is admissible in the sense that every internal label \underline{l} belongs to $\langle l^1, l^2 \rangle$, where l^1 and l^2 label the starting points of the incoming lines, corresponds to a unique such final subspace. The binary tree of the reduction scheme used here is



and admissible internal labelings are given by the sequences $\underline{l} \in R(\underline{j})$. \blacklozenge

To define the isomorphism $\varphi_{\underline{j}}$, we choose ladder bases in the irreducible subspaces $H_{\underline{j}, \underline{l}}$. Denote their elements by $|\underline{j}, \underline{l}, m\rangle$, where $m = -l^N, -l^N + 1, \dots, l^N$. Then,

$$\{|\underline{j}, \underline{l}, m\rangle : \underline{l} \in R(\underline{j}), m = -l^N, -l^N + 1, \dots, l^N\}$$

is an orthonormal basis in $H_{\underline{j}}$. Any other choice of ladder basis would yield the same basis vectors but multiplied by a phase which depends on \underline{l} only. For given $\underline{j} \in \langle \underline{j} \rangle$, we can use the sequences $\underline{l} \in R(\underline{j}, \underline{j})$ to label the copies of $H_{\underline{j}}$ in the direct sum decomposition of the target space of $\varphi_{\underline{j}}$. As a consequence, the natural projections and injections related with this decomposition read $\text{pr}_{\underline{j}, \underline{l}}^{\underline{j}}$ and $\text{i}_{\underline{j}, \underline{l}}^{\underline{j}}$, respectively, the basis functions read $(\chi_{\underline{j}})_{\underline{l}, \underline{l}'}^{\underline{j}}$, and the endomorphisms appearing in their definition read $A_{\underline{l}, \underline{l}'}^{jj}$. We define $\varphi_{\underline{j}}$ by

$$\varphi_{\underline{j}}(|\underline{j}, \underline{l}, m\rangle) := \text{i}_{\underline{j}, \underline{l}}^{\underline{j}}(|\underline{j}, m\rangle),$$

where $|\underline{j}, m\rangle$ denotes the elements of the orthonormal ladder basis in $H_{\underline{j}}$. Using (3.12) and the relation $\text{pr}_{\underline{j}, \underline{l}}^{\underline{j}} \circ \text{i}_{\underline{j}, \underline{l}'}^{\underline{j}} = \delta_{\underline{l}, \underline{l}'} \text{id}_{H_{\underline{j}}}$, for $\underline{j} \in \langle \underline{j} \rangle$, $\underline{l}, \underline{l}' \in R(\underline{j}, \underline{j})$, and $\underline{l}'' \in R(\underline{j})$, we compute

$$A_{\underline{l}, \underline{l}'}^{jj}(|\underline{j}, \underline{l}'', m\rangle) = \frac{1}{\sqrt{d_{\underline{j}}}} (\varphi_{\underline{j}})^{-1} \circ \text{i}_{\underline{j}, \underline{l}}^{\underline{j}} \circ \text{pr}_{\underline{j}, \underline{l}'}^{\underline{j}} \circ \varphi_{\underline{j}}(|\underline{j}, \underline{l}'', m\rangle) = \frac{\delta_{\underline{l}', \underline{l}''}}{\sqrt{d_{\underline{j}}}} |\underline{j}, \underline{l}, m\rangle.$$

This implies

$$A_{\underline{l}, \underline{l}'}^{jj} = \frac{1}{\sqrt{d_{\underline{j}}}} \sum_{m=-j}^j |\underline{j}, \underline{l}, m\rangle \langle \underline{j}, \underline{l}', m|, \quad \underline{l}, \underline{l}' \in R(\underline{j}, \underline{j}), \quad (4.3)$$

and

$$(\chi_{\underline{j}})_{\underline{l}, \underline{l}'}^{\underline{j}}(a) = \sqrt{\frac{d_{\underline{j}}}{d_{\underline{j}}}} \sum_{m=-j}^j \langle \underline{j}, \underline{l}', m | D^{\underline{j}}(a) | \underline{j}, \underline{l}, m \rangle, \quad \underline{l}, \underline{l}' \in R(\underline{j}, \underline{j}). \quad (4.4)$$

For later use, we express these functions in terms of the matrix entry functions $D_{m_i, m_i'}^i$, $i = 1, \dots, N$. For spins s_1, s_2, s and spin projections m_1, m_2, m , let

$$C_{m_1, m_2, m}^{s_1, s_2, s} := \langle s_1, m_1 | \otimes \langle s_2, m_2 | | s_1, s_2; s, m \rangle$$

denote the Clebsch-Gordan coefficients. Here, $|s_1, s_2; s, m\rangle$ denote the elements of the ladder basis in the irreducible subspace of spin s in $H_{s_1} \otimes H_{s_2}$ whenever $s \in \langle s_1, s_2 \rangle$ and the zero vector otherwise.

Proposition 4.2. We have

$$(\chi_{\underline{j}, \underline{l}, \underline{l}'}^j(\underline{a}) = \sqrt{\frac{d_j}{d_j}} \sum_{m=-j}^j \sum_{\underline{m}} \sum_{\underline{m}'} C(\underline{j}, \underline{l}, \underline{m}) C(\underline{j}, \underline{l}', \underline{m}') D_{m_1', m_1}^{j_1} (a_1) \cdots D_{m_N', m_N}^{j_N} (a_N),$$

where $\sum_{\underline{m}}$ means the sum over all sequences $\underline{m} = (m_1, \dots, m_N)$ such that

$$m_i = -j^i, \dots, j^i \text{ for } i = 1, \dots, N, \quad m_1 + \dots + m_N = m,$$

and where

$$C(\underline{j}, \underline{l}, \underline{m}) = C_{m_1, m_2, m_1+m_2}^{j^1, j^2, l^2} C_{m_1+m_2, m_3, m_1+m_2+m_3}^{l^2, j^3, l^3} \cdots C_{m_1+\dots+m_{N-1}, m_N, m}^{l^{N-1}, j^N, l^N}.$$

Proof. Using the tensor basis in H_j , given by the vectors

$$|\underline{j}, \underline{m}\rangle := |j_1, m_1\rangle \otimes \cdots \otimes |j_N, m_N\rangle, \quad m_i = -j_i, \dots, j_i, \quad i = 1, \dots, N,$$

formula (4.4) can be rewritten as

$$(\chi_{\underline{j}, \underline{l}, \underline{l}'}^j(\underline{a}) = \sqrt{\frac{d_j}{d_j}} \sum_{m=-j}^j \sum_{\underline{m}} \sum_{\underline{m}'} \langle \underline{j}, \underline{l}', \underline{m} | \underline{j}, \underline{m} | \underline{j}, \underline{l}, \underline{m} \rangle D_{m_1', m_1}^{j_1} (a_1) \cdots D_{m_N', m_N}^{j_N} (a_N).$$

To compute the scalar products, we expand $|\underline{j}, \underline{l}, \underline{m}\rangle$ with respect to $|\underline{j}, \underline{m}\rangle$. Denote $\underline{j} := (j^1, \dots, j^{N-1})$ and $\underline{l} := (l^1, \dots, l^{N-1})$ and consider the irreducible subspace $H_{\underline{j}, \underline{l}}$ of $H_{\underline{j}}$ with its ladder basis $\{|\underline{j}, \underline{l}, \underline{m}\rangle : m = -l^{N-1}, \dots, l^{N-1}\}$. By construction, $H_{\underline{j}, \underline{l}}$ is the irreducible subspace of spin l^N in $H_{\underline{j}, \underline{l}} \otimes H_{j^N}$ and $|\underline{j}, \underline{l}, \underline{m}\rangle$ are the elements of the ladder basis in that subspace. Hence,

$$|\underline{j}, \underline{l}, \underline{m}\rangle = \sum_{n_N = -l^{N-1}}^{l^{N-1}} \sum_{m_N = -j^N}^{j^N} C_{n_N, m_N, m}^{l^{N-1}, j^N, l^N} |\underline{j}, \underline{l}, n_N\rangle \otimes |j^N, m_N\rangle.$$

Iterating this argument, we find that the expansion of $|\underline{j}, \underline{l}, \underline{m}\rangle$ is given by

$$\sum_{m_i, n_i} C_{n_N, m_N, m}^{l^{N-1}, j^N, l^N} C_{n_{N-1}, m_{N-1}, n_N}^{l^{N-2}, j^{N-1}, l^{N-1}} \cdots C_{n_2, m_2, n_3}^{l^1, j^2, l^2} |j^1, n_2\rangle \otimes |j^2, m_2\rangle \otimes \cdots \otimes |j^N, m_N\rangle,$$

where the sum runs over $n_i = -l^{i-1}, \dots, l^{i-1}$ and $m_i = -j^i, \dots, j^i$ for $i = 2, \dots, N$. Putting $m_1 = n_2$ and taking into account that the Clebsch-Gordan coefficients vanish unless the first two spin projections add up to the third one, we find that in the sum over n_3, \dots, n_N , only the terms with

$$n_3 = m_1 + m_2, \quad n_4 = m_1 + \dots + m_3, \quad \dots, \quad n_N = m_1 + \dots + m_{N-1}$$

survive. As a result, we obtain

$$|\underline{j}, \underline{l}, \underline{m}\rangle = \sum_{\underline{m}} C(\underline{j}, \underline{l}, \underline{m}) |\underline{j}, \underline{m}\rangle.$$

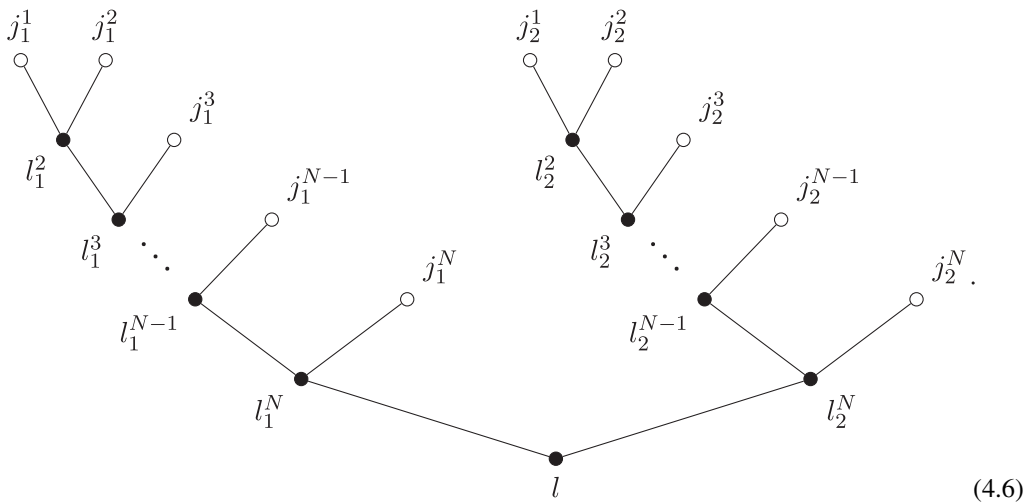
Plugging this into the above formula for $(\chi_{\underline{j}, \underline{l}, \underline{l}'}^j(\underline{a}))$ and taking into account that the Clebsch-Gordan coefficients are real, we obtain the assertion. \square

Next, we compute the coefficients U in the multiplication law for the basis functions given by Proposition 3.10. For this purpose, we have to determine the unitary G -representation isomorphisms Φ_{j_1, j_2} and Ψ_{j_1, j_2} introduced in Subsection III D.

First, consider $\Phi_{j_1 j_2}$. Recall that $\Phi_{j_1 j_2} = \phi_{j_1 j_2} \circ (\varphi_{j_1} \otimes \varphi_{j_2})$. Here, $\phi_{j_1 j_2}$ is given by a unitary G -representation isomorphism

$$H_{j_1} \otimes H_{j_2} \rightarrow \bigoplus_{j \in \langle j_1 j_2 \rangle} H_j \quad (4.5)$$

for every pair j_1, j_2 with $j_1 \in \langle j_1 \rangle$ and $j_2 \in \langle j_2 \rangle$. Since the multiplicities are 1 here, we may omit the corresponding index in our notation. Another consequence is that the isomorphism (4.5) is determined up to a phase on every H_j . We choose these phases in accordance with the standard choice of the Clebsch-Gordan coefficients so that (4.5) is given by these coefficients. Then, $\Phi_{j_1 j_2}$ is uniquely determined by the choice of φ_j for every j and hence by the choice of the reduction scheme for N -fold tensor products of $SU(2)$ -irreps. To write it down explicitly, we decompose $H_{j_1} \otimes H_{j_2}$ into irreducible subspaces according to the following reduction scheme:



This leads to irreducible subspaces labeled by $l_1 \in R(j_1)$, $l_2 \in R(j_2)$, and $l \in \langle l_1^N, l_2^N \rangle$. In each subspace, we choose an orthonormal ladder basis and denote its elements by $|j_1, j_2; l_1, l_2; l, m\rangle$, $m = -l, \dots, l$. Then,

$$\{|j_1, j_2; l_1, l_2; l, m\rangle : l_1 \in R(j_1), l_2 \in R(j_2), l \in \langle l_1^N, l_2^N \rangle, m = -l, \dots, l\}$$

is an orthonormal basis in $H_{j_1} \otimes H_{j_2}$, and $\Phi_{j_1 j_2}$ is given by

$$\Phi_{j_1 j_2}(|j_1, j_2; l_1, l_2; l, m\rangle) = i_{l_1^N, l_2^N, l_1, l_2, l}^{j_1, j_2}(|l, m\rangle),$$

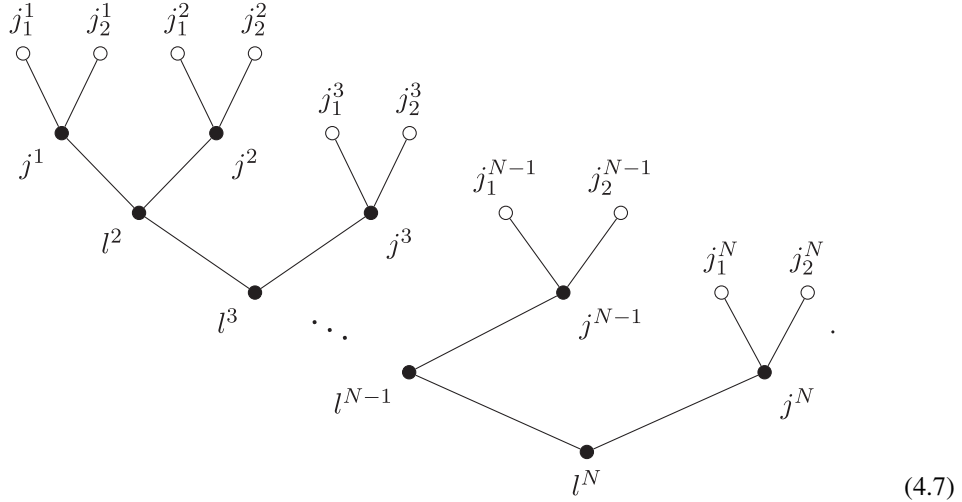
where $i_{l_1^N, l_2^N, l_1, l_2, l}^{j_1, j_2}$ denotes the natural injection associated with the decomposition (3.18) (here, by our specific choice of notation, the labels l_1^N and l_2^N are actually redundant).

Now, consider the unitary G -representation isomorphism $\Psi_{j_1 j_2}$. Denote $\langle j_1, j_2 \rangle = \prod_{i=1}^N \langle j_1^i, j_2^i \rangle$. Recall that $\Psi_{j_1 j_2} = \psi_{j_1 j_2}^2 \circ \psi_{j_1 j_2}^1$, where

$$\psi_{j_1 j_2}^1 : H_{j_1} \otimes H_{j_2} \rightarrow \bigoplus_{\underline{j} \in \langle j_1, j_2 \rangle} H_{\underline{j}}$$

is a unitary G^N -representation isomorphism and $\psi_{j_1 j_2}^2$ acts on every summand $H_{\underline{j}}$ as $\varphi_{\underline{j}}$. Since for every factor of G^N , $\psi_{j_1 j_2}^1$ boils down to an isomorphism of the type (4.5), the multiplicities are 1 as well and so we may omit the corresponding index in our notation. This also implies that $\psi_{j_1 j_2}^1$

is unique up to a phase for every factor of G^N and the corresponding irreducible factor of $H_{\underline{j}}$. As before, we choose these phases so that $\psi_{j_1 j_2}^1$ is given by the appropriate Clebsch-Gordan coefficients. Then, $\Psi_{j_1 j_2}$, like $\Phi_{j_1 j_2}$, is uniquely determined by the choice of the reduction scheme for N -fold tensor products of $SU(2)$ -irreps. To write it down explicitly, we decompose $H_{j_1} \otimes H_{j_2}$ into irreducible subspaces according to the following reduction scheme:



This leads to irreducible subspaces labeled by $\underline{j} \in \langle j_1, j_2 \rangle$ and $\underline{l} \in R(\underline{j})$. In each subspace, we choose an orthonormal ladder basis and denote its elements by $|\underline{j}_1, \underline{j}_2; \underline{j}, \underline{l}, m\rangle$, $m = -l^N, \dots, l^N$. Then,

$$\{|\underline{j}_1, \underline{j}_2; \underline{j}, \underline{l}, m\rangle : \underline{j} \in \langle j_1, j_2 \rangle, \underline{l} \in R(\underline{j}), m = -l^N, \dots, l^N\}$$

is an orthonormal basis in $H_{j_1} \otimes H_{j_2}$, and $\Psi_{j_1 j_2}$ is given by

$$\Psi_{j_1 j_2}(|\underline{j}_1, \underline{j}_2; \underline{j}, \underline{l}, m\rangle) = \underset{j, l^N, \underline{l}}{i_{j_1 j_2}^{j_1 j_2}}(|l^N, m\rangle),$$

where $\underset{j, l^N, \underline{l}}{i_{j_1 j_2}^{j_1 j_2}}$ denotes the natural injection associated with the decomposition (3.22) (where, by our specific choice of notation, the label l^N is redundant).

Proposition 4.3. *In the case of $G = SU(2)$, the multiplication law for the basis functions $(\chi_{\underline{j}}^{\mathbb{C}})_{\underline{l}, \underline{l}'}^j$ reads*

$$(\chi_{j_1}^{\mathbb{C}})_{l_1, l_1'}^{j_1} \cdot (\chi_{j_2}^{\mathbb{C}})_{l_2, l_2'}^{j_2} = \sqrt{\frac{d_{j_1} d_{j_2}}{d_{j_1} d_{j_2}}} \sum_{\underline{j} \in \langle j_1, j_2 \rangle} \sum_{\underline{j} \in \langle j_1, j_2 \rangle} \sum_{\underline{l}, \underline{l}' \in R(\underline{j})} \sqrt{\frac{d_{\underline{j}}}{d_{\underline{j}}}} U_{j_1 j_2}(\underline{j}, \underline{l}; l_1, l_2) U_{j_1 j_2}(\underline{j}, \underline{l}'; l_1', l_2') (\chi_{\underline{j}}^{\mathbb{C}})_{\underline{l}, \underline{l}'}^j,$$

where

$$U_{j_1 j_2}(\underline{j}, \underline{l}; l_1, l_2) = \langle \underline{j}_1, \underline{j}_2; \underline{j}, \underline{l}, m | \underline{j}_1, \underline{j}_2; l_1, l_2; j, m \rangle \quad (4.8)$$

for every $\underline{j} \in \langle j_1, j_2 \rangle$, $\underline{j} \in \langle j_1, j_2 \rangle$, and $\underline{l} \in R(\underline{j})$ and for any admissible m .

The coefficients

$$U_{j_1 j_2}(\underline{j}, \underline{l}; l_1, l_2) = \langle \underline{j}_1, \underline{j}_2; \underline{j}, \underline{l}, m | \underline{j}_1, \underline{j}_2; l_1, l_2; j, m \rangle$$

are the recoupling coefficients for the reduction schemes (4.6) and (4.7). Up to normalization, they are given by what is known as $3(2N-1)j$ symbols; see Topic 12 in Ref. 5 for details.

Proof. By Proposition 3.10, it suffices to compute the coefficients $U_{j_1 j_2}(\underline{j}, \underline{l}; l_1, l_2)$. According to (3.25), they are defined by

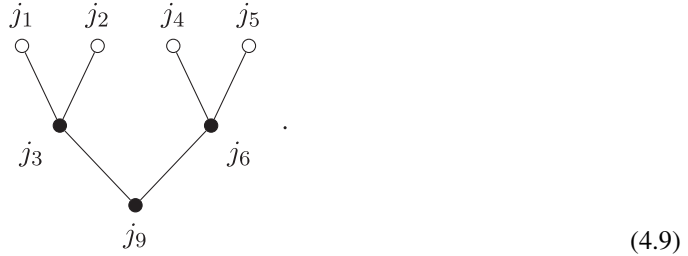
$$\text{pr}_{j_1 j_2}^{j_1 j_2} \circ \Psi_{j_1 j_2} \circ \Phi_{j_1 j_2}^{-1} \circ \Gamma_{j_1 j_2, l_1, l_2, j}^{j_1 j_2} = U_{j_1 j_2}(\underline{j}, \underline{l}; l_1, l_2) \text{id}_{H_j},$$

where $j \in \langle j_1, j_2 \rangle$, $j \in \langle j_1, j_2 \rangle$, and $\underline{l} \in R(j, j)$. Evaluating the left-hand side on a ladder basis vector $|j, m\rangle$ and plugging in a unit, we calculate

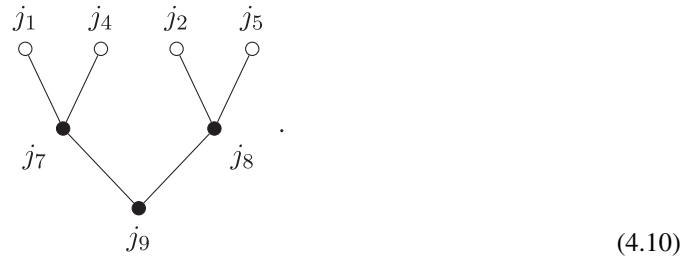
$$\begin{aligned} & \text{pr}_{j_1 j_2}^{j_1 j_2} \circ \Psi_{j_1 j_2} \circ \Phi_{j_1 j_2}^{-1} \circ \Gamma_{j_1 j_2, l_1, l_2, j}^{j_1 j_2}(|j, m\rangle) \\ &= \text{pr}_{j_1 j_2}^{j_1 j_2} \circ \Psi_{j_1 j_2}(|j_1, j_2; \underline{l}_1, \underline{l}_2; j, m\rangle) \\ &= \sum_{j', l'} \langle j_1, j_2; j', l', m | j_1, j_2; \underline{l}_1, \underline{l}_2; j, m \rangle \text{pr}_{j_1 j_2}^{j_1 j_2} \circ \Psi_{j_1 j_2}(|j_1, j_2; j', l', m\rangle) \\ &= \sum_{j', l'} \langle j_1, j_2; j', l', m | j_1, j_2; \underline{l}_1, \underline{l}_2; j, m \rangle \text{pr}_{j_1 j_2}^{j_1 j_2} \circ \Gamma_{j', l', j}^{j_1 j_2}(|l'^N, m\rangle) \\ &= \langle j_1, j_2; \underline{j}, \underline{l}, m | j_1, j_2; \underline{l}_1, \underline{l}_2; j, m \rangle |j, m\rangle. \end{aligned}$$

This yields (4.8). The multiplication law follows then by observing that the coefficients $U_{j_1 j_2}(\underline{j}, \underline{l}; l_1, l_2)$ are real. \square

The recoupling coefficients $U_{j_1 j_2}(\underline{j}, \underline{l}; l_1, l_2)$ can be expressed in terms of the recoupling coefficients for $N = 2$, that is, for a tensor product of four $\text{SU}(2)$ -irreps as follows. Given four spins j_1, j_2, j_4, j_5 , the tensor product $D^{j_1} \otimes D^{j_2} \otimes D^{j_4} \otimes D^{j_5}$ can be decomposed, on the one hand, into irreducible subspaces labeled by j_3, j_6, j_9 according to the following reduction scheme:



Let $\{|(j_1, j_2), (j_4, j_5); j_3, j_6; j_9, m\rangle : m = -j_9, \dots, j_9\}$ be the ladder bases in these subspaces chosen in accordance with the definition of the Clebsch-Gordan coefficients. (The notation for the basis elements corresponds to the general notation used here, specified to $N = 2$; the conventional notation is $|((j_1 j_2) j_3, (j_4 j_5) j_6) j_9, m\rangle$.) On the other hand, this tensor product can be decomposed into irreducible subspaces labeled by j_7, j_8, j_9 according to the following reduction scheme:



Let $\{|(j_1, j_2), (j_4, j_5); (j_7, j_8), j_9, m\rangle : m = -j_9, \dots, j_9\}$ be the ladder bases in these subspaces, again chosen in accordance with the definition of the Clebsch-Gordan coefficients (i.e., $|((j_1 j_4) j_7, (j_2 j_5) j_8) j_9, m\rangle$ in the conventional notation). It is common to denote the recoupling coefficients between these two reduction schemes by

$$\begin{pmatrix} j_1 & j_2 & j_3 \\ j_4 & j_5 & j_6 \\ j_7 & j_8 & j_9 \end{pmatrix} := \langle (j_1, j_2), (j_4, j_5); j_3, j_6; j_9, m | (j_1, j_2), (j_4, j_5); (j_7, j_8), j_9, m \rangle$$

(the right-hand side does not depend on m). These coefficients are related with Wigner's $9j$ symbols via

$$\begin{pmatrix} j_1 & j_2 & j_3 \\ j_4 & j_5 & j_6 \\ j_7 & j_8 & j_9 \end{pmatrix} = \sqrt{(2j_3+1)(2j_6+1)(2j_7+1)(2j_8+1)} \begin{Bmatrix} j_1 & j_2 & j_3 \\ j_4 & j_5 & j_6 \\ j_7 & j_8 & j_9 \end{Bmatrix}.$$

Lemma 4.4. *The recoupling coefficients $U_{j_1 j_2}(j, l; l_1, l_2)$ are given by*

$$U_{j_1 j_2}(j, l; l_1, l_2) = \prod_{i=2}^N \begin{pmatrix} l_1^{i-1} & l_2^{i-1} & l^{i-1} \\ j_1^i & j_2^i & j^i \\ l_1^i & l_2^i & l^i \end{pmatrix},$$

where $l_1^1 := j_1^1$, $l_2^1 := j_2^1$, and $l^1 := j^1$.

Proof. See [Appendix B](#). □

For general recursion formulae for $3nj$ symbols, see Ref. [41](#). Since there exist efficient calculators for $9j$ symbols, provided, e.g., by the Python library SymPy²⁹ or online by Anthony Stone's Wigner coefficient calculator,³⁸ Lemma 4.4 provides an explicit knowledge of the multiplication law in the commutative algebra \mathcal{R} for $SU(2)^N$.

B. The orbit type strata

Recall that for $G = SU(2)$, we have $G_{\mathbb{C}} = SL(2, \mathbb{C})$, $\mathfrak{g} = \mathfrak{su}(2)$ and $\mathfrak{g}_{\mathbb{C}} = \mathfrak{sl}(2, \mathbb{C})$. For convenience, we keep the notation G and $G_{\mathbb{C}}$. Let Z denote the center of G . Clearly, this is also the center of $G_{\mathbb{C}}$. Let $T \subset G$ denote the subgroup of diagonal matrices, and let \mathfrak{t} be its Lie algebra. Clearly, T is a maximal toral subgroup isomorphic to $U(1)$.

Let us briefly recall the orbit type strata \mathcal{P}_{τ} of \mathcal{P} in terms of subsets $(G_{\mathbb{C}}^N)_{\tau}$ of $G_{\mathbb{C}}^N$. For details, see Ref. [11](#). First, one determines the orbit types of the lifted action of G on T^*G^N . There are three of them and these can be labeled by G , T , and Z , where Z is the principal orbit type. The corresponding orbit type subsets of $G^N \times \mathfrak{g}^N$ are as follows.

(G) An element $(\underline{a}, \underline{A}) \in G^N \times \mathfrak{g}^N$ has orbit type G if and only if

$$(\underline{a}, \underline{A}) \in Z^N \times \{0\}^N.$$

(T) An element $(\underline{a}, \underline{A}) \in G^N \times \mathfrak{g}^N$ has orbit type T if and only if it is conjugate to an element of the subset

$$(T^N \times \mathfrak{t}^N) \setminus (Z^N \times \{0\}^N).$$

Since conjugation by an element of G commutes with taking commutators, for every element $(\underline{a}, \underline{A})$ of orbit type T , the entries $(a_1, \dots, a_N, A_1, \dots, A_N)$ commute pairwise. Conversely, if for an element $(\underline{a}, \underline{A})$ all its entries commute pairwise, then they are simultaneously diagonalizable and hence they belong to the orbit type T .

(Z) An element $(\underline{a}, \underline{A}) \in G^N \times \mathfrak{g}^N$ has orbit type Z if and only if it does not have orbit type T or G , i.e., if and only if it is not conjugate to an element of $T^N \times \mathfrak{t}^N$, that is, if and only if not all entries of $(\underline{a}, \underline{A})$ commute pairwise.

Next, one intersects the orbit type subsets with the momentum level set $\mu^{-1}(0)$, takes the quotient of $\mu^{-1}(0)$ with respect to the G -action, and passes to connected components. This yields the following.

(G) There exist 2^N orbit type strata of orbit type G , each of which consists of a single point representing the (trivial) orbit of an element of $Z^N \times \{0\}^N$. Since such an element is of the

form $(\nu_1 \mathbb{1}, \dots, \nu_N \mathbb{1}, 0, \dots, 0)$ for some sequence of signs $\underline{\nu} = (\nu_1, \dots, \nu_N)$, we denote the corresponding stratum by $\mathcal{P}_{\underline{\nu}}$.

- (T) Since $Z^N \times \{0\}^N$ consists of finitely many points and $T^N \times \mathfrak{t}^N$ has dimension at least 2, the complement $(T^N \times \mathfrak{t}^N) \setminus (Z^N \times \{0\}^N)$ is connected. Since the subset of \mathcal{P} of orbit type T is the image of the subset

$$(T^N \times \mathfrak{t}^N) \setminus (Z^N \times \{0\}^N) \subset \mu^{-1}(0)$$

under the natural projection $\mu^{-1}(0) \rightarrow \mathcal{P}$, it is connected, too. Hence, it forms an orbit type stratum. We denote this stratum by \mathcal{P}_T .

- (Z) Since \mathfrak{g}^* has dimension 3, the level set $\mu^{-1}(0)$ generically has dimension $2N \cdot 3 - 3 = 3(2N - 1)$. On the other hand, since T has dimension 1 and the elements of $T^N \times \mathfrak{t}^N$ have stabilizer T under the action of G , the subset of $G^N \times \mathfrak{g}^N$ of orbit type T has dimension $2N \cdot 1 + (3 - 1) = 2(N + 1)$. Hence, if the orbit type Z occurs in \mathcal{P} , i.e., if $N \geq 2$, then the subset of $\mu^{-1}(0)$ generated from $T^N \times \mathfrak{t}^N$ by the action of G has codimension

$$3(2N - 1) - 2(N + 1) = 4N - 5 \geq 3.$$

Therefore, its complement is connected. Since the complement coincides with the subset of $\mu^{-1}(0)$ of orbit type Z , the subset of \mathcal{P} of this orbit type is connected. Hence, it forms an orbit type stratum. We denote this stratum by \mathcal{P}_Z .

One can visualize the set of strata and their partial ordering by a Hasse diagram; see Ref. 11. Finally, one transports the above results to $G_{\mathbb{C}}^N$, that is, for each of the above strata, one finds the subset $(G_{\mathbb{C}}^N)_{\tau}$ of $G_{\mathbb{C}}^N$. It suffices to do this for every sequence of signs $\underline{\nu} = (\nu_1, \dots, \nu_N)$ and for T . Let $T_{\mathbb{C}} \subset G_{\mathbb{C}}$ denote the subgroup of diagonal matrices. One obtains the following.

Theorem 4.5. *Let $\underline{a} \in G_{\mathbb{C}}^N$. Then,*

1. $\underline{a} \in (G_{\mathbb{C}}^N)_{\underline{\nu}}$ if and only if \underline{a} is orbit closure equivalent to $(\nu_1 \mathbb{1}, \dots, \nu_N \mathbb{1})$,
2. $\underline{a} \in (G_{\mathbb{C}}^N)_T$ if and only if \underline{a} is orbit closure equivalent to an element of $T_{\mathbb{C}}^N \setminus Z^N$.

C. Zero locus and radical ideal conditions

In this subsection, for the strata τ found above, we recall from Ref. 11 the finite subsets R_{τ} of \mathcal{R} having the corresponding orbit type subset $(G_{\mathbb{C}}^N)_{\tau}$ as their zero locus and satisfying the radical ideal condition. Since $\tau = Z$ corresponds to the principal stratum and hence $\mathcal{H}_Z = \mathcal{H}$, it suffices to discuss the secondary strata $\tau = \underline{\nu}$ and $\tau = T$. For this purpose, we define the following G -invariant representative functions:

$$\begin{aligned} p_{rs}^T(\underline{a}) &:= \text{tr}([a_r, a_s]^2), & 1 \leq r < s \leq N, \\ p_{rst}^T(\underline{a}) &:= \text{tr}([a_r, a_s]a_t), & 1 \leq r < s < t \leq N. \end{aligned}$$

For the strata labeled by sequences of signs $\underline{\nu}$, one obtains the following.

Theorem 4.6. *The subset $(G_{\mathbb{C}}^N)_{\underline{\nu}} \subset G_{\mathbb{C}}^N$ is the set of common zeros of the $G_{\mathbb{C}}$ -invariant functions p_{rs}^T with $1 \leq r < s \leq N$, p_{rst}^T with $1 \leq r < s < t \leq N$, and*

$$p_r^{\underline{\nu}}(\underline{a}) := \text{tr}(a_r) - \nu_r 2, \quad r = 1, \dots, N.$$

Remark 4.7. Instead of using Theorem 4.6, one can construct the subspace $\mathcal{H}_{\underline{\nu}}$ associated with the stratum $\mathcal{P}_{\underline{\nu}}$ directly as follows. Let $\{\psi_{\alpha} : \alpha \in A\}$ be an orthonormal basis of \mathcal{H} which contains a constant function ψ_0 . Clearly, the basis provided by Proposition 3.7 is of that type. Since for a continuous invariant function ψ , the condition to vanish on $(G_{\mathbb{C}}^N)_{\underline{\nu}}$ is equivalent to the condition

$\psi(v_1 \mathbb{1}, \dots, v_N \mathbb{1}) = 0$, the vanishing subspace $\mathcal{V}_{\underline{v}}$ of the stratum $\mathcal{P}_{\underline{v}}$, given by (3.7), is spanned by the elements

$$\psi_\alpha - \psi_\alpha(v_1 \mathbb{1}, \dots, v_N \mathbb{1}) \mathbb{1}, \quad \alpha \in A, \alpha \neq 0,$$

where $\mathbb{1}$ denotes the constant function with value 1. One proves that $\mathcal{H}_{\underline{v}}$ is spanned by the single element

$$\psi_{\underline{v}} = \frac{1}{C_{\underline{v}}} \sum_{\beta \in A} \overline{\psi_\beta(v_1 \mathbb{1}, \dots, v_N \mathbb{1})} \psi_\beta,$$

where $C_{\underline{v}}$ is a normalization constant. See Remark 5.4 in Ref. 11 for the details. \square

By this remark, checking the radical ideal condition is relevant for the stratum \mathcal{P}_T only.

Theorem 4.8. *The topological closure $\overline{(G_{\mathbb{C}}^N)_T}$ is the set of common zeros of the G -invariant representative functions*

$$p_{rs}^T, \quad 1 \leq r < s \leq N, \quad p_{rst}^T, \quad 1 \leq r < s < t \leq N. \quad (4.11)$$

The ideal in \mathcal{R} generated by these functions is a radical ideal.

The proof of the radical ideal condition is the hard part of Ref. 11.

For the construction of the subspace \mathcal{H}_T associated with the stratum T , it will be convenient to express the functions p_{rs}^T and p_{rst}^T in terms of the basis functions $(\chi_{\underline{j}}^{\mathbb{C}})_{\underline{l}, \underline{l}'}$ introduced in Sec. IV A. It will turn out that the functions p_{rs}^T are linear combinations of basis functions with \underline{j} having entries j_r at r , j_s at s , and 0 elsewhere. For such a sequence, we write $\underline{j} = (j_r r, j_s s)$. The corresponding sequences $\underline{l} \in R(\underline{j})$ have entries $l^1 = \dots = l^{r-1} = 0$, $l^r = \dots = l^{s-1} = j_r$, and $l^s = \dots = l^N = j_s$, where $j \in \langle j_r, j_s \rangle$. Hence, for given \underline{j} , there is only one sequence \underline{l} in $R(\underline{j}, \underline{j})$ so that we may omit the labels \underline{l} and \underline{l}' in the notation. In a similar way, the functions p_{rst}^T will turn out to be linear combinations of basis functions with \underline{j} having entries j_r at r , j_s at s , j_t at t , and 0 elsewhere. For such a sequence, we write $\underline{j} = (j_r r, j_s s, j_t t)$. Here, the sequences $\underline{l} \in R(\underline{j})$ have entries $l^1 = \dots = l^{r-1} = 0$, $l^r = \dots = l^{s-1} = j_r$, $l^s = \dots = l^{t-1} = j_s$, and $l^t = \dots = l^N = j_t$, where $l \in \langle j_r, j_s \rangle$ and $j \in \langle l, j_s \rangle$. That is, they are labeled by a single intermediate spin l so that in our notation we may replace the labels \underline{l} and \underline{l}' by l and l' , respectively.

Lemma 4.9. *The functions p_{rs}^T and p_{rst}^T on $G_{\mathbb{C}}^N$ are given by*

$$p_{rs}^T = (\chi_{(1r, 0s)}^{\mathbb{C}})^1 + (\chi_{(0r, 1s)}^{\mathbb{C}})^1 + (\chi_{(1r, 1s)}^{\mathbb{C}})^0 - \frac{2}{\sqrt{3}} (\chi_{(1r, 1s)}^{\mathbb{C}})^1 - 3, \quad (4.12)$$

$$p_{rst}^T = \frac{\sqrt{3}}{2} \left((\chi_{(\frac{1}{2}r, \frac{1}{2}s, \frac{1}{2}t)}^{\mathbb{C}})_{0,1}^{\frac{1}{2}} - (\chi_{(\frac{1}{2}r, \frac{1}{2}s, \frac{1}{2}t)}^{\mathbb{C}})_{1,0}^{\frac{1}{2}} \right). \quad (4.13)$$

Proof. According to the last statement of Corollary 3.8, the expansion coefficients of p_{rs}^T and p_{rst}^T with respect to the basis $\{(\chi_{\underline{j}}^{\mathbb{C}})_{\underline{l}, \underline{l}'}^j\}$ in \mathcal{H} coincide with the expansion coefficients of their restrictions to G^N with respect to the basis $\{(\chi_{\underline{j}}^{\mathbb{C}})_{\underline{l}, \underline{l}'}^j\}$ in $L^2(G^N)^G$. Hence, it suffices to determine the latter. By an abuse of notation, in what follows, p_{rs}^T and p_{rst}^T mean the restrictions to G^N .

First, consider p_{rs}^T . We have

$$p_{rs}^T = 2 \operatorname{tr}((a_r a_s)^2) - 2 \operatorname{tr}(a_r^2 a_s^2).$$

For the second term, we use $a = D^{\frac{1}{2}}(a)$ and $D^{\frac{1}{2}} \otimes D^{\frac{1}{2}} = D^0 \oplus D^1$ to calculate

$$\begin{aligned}
\text{tr}(a_r^2 a_s^2) &= \sum_{m_i = \pm \frac{1}{2}} D_{m_1 m_2}^{\frac{1}{2}}(a_r) D_{m_2 m_3}^{\frac{1}{2}}(a_r) D_{m_3 m_4}^{\frac{1}{2}}(a_s) D_{m_4 m_1}^{\frac{1}{2}}(a_s) \\
&= \sum_{m_i = \pm \frac{1}{2}} \left\langle \left\langle \frac{1}{2}, m_1 \right| \otimes \left\langle \frac{1}{2}, m_2 \right| \left| D^{\frac{1}{2}}(a_r) \otimes D^{\frac{1}{2}}(a_r) \right| \left| \frac{1}{2}, m_2 \right\rangle \otimes \left| \frac{1}{2}, m_3 \right\rangle \right\rangle \\
&\quad \cdot \left\langle \left\langle \frac{1}{2}, m_3 \right| \otimes \left\langle \frac{1}{2}, m_4 \right| \left| D^{\frac{1}{2}}(a_s) \otimes D^{\frac{1}{2}}(a_s) \right| \left| \frac{1}{2}, m_4 \right\rangle \otimes \left| \frac{1}{2}, m_1 \right\rangle \right\rangle \\
&= \sum_{m_i = \pm \frac{1}{2}} \sum_{j_r, j_s=0}^1 \sum_{n_r, n'_r=-j_r}^{j_r} \sum_{n_s, n'_s=-j_s}^{j_s} \\
&\quad C_{m_1 m_2 n_r}^{\frac{1}{2} \frac{1}{2} j_r} C_{m_2 m_3 n'_r}^{\frac{1}{2} \frac{1}{2} j_r} C_{m_3 m_4 n'_s}^{\frac{1}{2} \frac{1}{2} j_s} C_{m_4 m_1 n_s}^{\frac{1}{2} \frac{1}{2} j_s} D_{n_r n'_r}^{j_r}(a_r) D_{n_s n'_s}^{j_s}(a_s).
\end{aligned}$$

Hence,

$$\begin{aligned}
\text{tr}(a_r^2 a_s^2) &= \sum_{j_r, j_s=0}^1 \sum_{n_r, n'_r=-j_r}^{j_r} \sum_{n_s, n'_s=-j_s}^{j_s} R_{n_r n'_r, n_s n'_s}^{j_r j_s} D_{n_r n'_r}^{j_r}(a_r) D_{n_s n'_s}^{j_s}(a_s), \\
R_{n_r n'_r, n_s n'_s}^{j_r j_s} &:= \sum_{m_i = \pm \frac{1}{2}} C_{m_1 m_2 n_r}^{\frac{1}{2} \frac{1}{2} j_r} C_{m_2 m_3 n'_r}^{\frac{1}{2} \frac{1}{2} j_r} C_{m_3 m_4 n'_s}^{\frac{1}{2} \frac{1}{2} j_s} C_{m_4 m_1 n_s}^{\frac{1}{2} \frac{1}{2} j_s}.
\end{aligned} \tag{4.14}$$

A similar calculation for the first term of p_{rs}^T , using the relations

$$C_{m_1 m_2 m}^{j_1 j_2 j} = (-1)^{j_1 + j_2 - j} C_{m_2 m_1 m}^{j_2 j_1 j}, \quad \sum_{m_1, m_2} C_{m_1 m_2 m}^{j_1 j_2 j} C_{m_1 m_2 m'}^{j_1 j_2 j'} = \delta_{jj'} \delta_{mm'},$$

yields

$$\begin{aligned}
\text{tr}((a_r a_s)^2) &= \sum_{j_r, j_s=0}^1 \sum_{n_r, n'_r=-j_r}^{j_r} \sum_{n_s, n'_s=-j_s}^{j_s} S_{n_r n'_r, n_s n'_s}^{j_r j_s} D_{n_r n'_r}^{j_r}(a_r) D_{n_s n'_s}^{j_s}(a_s), \\
S_{n_r n'_r, n_s n'_s}^{j_r j_s} &:= (-1)^{1-j_r} \delta_{j_r j_s} \delta_{n_r n'_r} \delta_{n'_r n_s}.
\end{aligned} \tag{4.15}$$

From (4.14) and (4.15), we conclude that p_{rs}^T is a linear combination of the basis functions $(\chi_{(j_r, r, j_s, s)})^j$ with $j_r, j_s = 0, 1$ and $j \in \langle j_r, j_s \rangle$. To compute the expansion coefficients, we use Proposition 4.2 to write

$$(\chi_{(j_r, r, j_s, s)})^j = \sqrt{\frac{d_{j_r} d_{j_s}}{d_j}} \sum_{m=-j}^j \sum_{\substack{n_r + n_s = m \\ n'_r + n'_s = m}} C_{n_r n_s m}^{j_r j_s j} C_{n'_r n'_s m}^{j_r j_s j} D_{n_r n'_r}^{j_r}(a_r) D_{n_s n'_s}^{j_s}(a_s) \tag{4.16}$$

and compute the scalar products $\langle (\chi_{(j_r, r, j_s, s)})^j | p_{rs}^T \rangle$ using the orthogonality relation

$$\langle D_{m_1 m_2}^j | D_{m'_1 m'_2}^{j'} \rangle = \frac{1}{d_j} \delta_{jj'} \delta_{m_1 m'_1} \delta_{m_2 m'_2}. \tag{4.17}$$

This results in

$$\langle (\chi_{(j_r, r, j_s, s)})^j | p_{rs}^T \rangle = \frac{2}{\sqrt{d_{j_r} d_{j_s} d_j}} \sum_{m=-j}^j \sum_{\substack{n_r + n_s = m \\ n'_r + n'_s = m}} C_{n_r n_s m}^{j_r j_s j} C_{n'_r n'_s m}^{j_r j_s j} (S_{n_r n'_r, n_s n'_s}^{j_r j_s} - R_{n_r n'_r, n_s n'_s}^{j_r j_s}).$$

Computation of the right-hand side for $j_r, j_s = 0, 1$ and $j \in \langle j_r, j_s \rangle$ yields (4.12). For computations involving products of Clebsch-Gordan coefficients, one may use, for example, the Clebsch-Gordan coefficient function of Mathematica⁴⁰ or the class `sympy.physics.quantum.cg.CG` provided by the Python library SymPy.²⁹

For p_{rst}^T , we proceed in an analogous way. Writing

$$p_{rst}^T(\underline{a}) = \sum_{m_i = \pm \frac{1}{2}} \left(D_{m_1 m_2}^{\frac{1}{2}}(a_r) D_{m_2 m_3}^{\frac{1}{2}}(a_s) - D_{m_1 m_2}^{\frac{1}{2}}(a_s) D_{m_2 m_3}^{\frac{1}{2}}(a_r) \right) D_{m_3 m_1}^{\frac{1}{2}}(a_t),$$

we see that p_{rst}^T is a linear combination of the basis functions $(\chi_{(\frac{1}{2}r, \frac{1}{2}s, \frac{1}{2}t)})_{ll'}^j$ with admissible j, l , and l' , i.e., with $j = \frac{1}{2}$ and $l, l' = 0, 1$ or with $j = \frac{3}{2}$ and $l = l' = 1$ (five functions altogether). According to Proposition 4.2,

$$\begin{aligned} (\chi_{(\frac{1}{2}r, \frac{1}{2}s, \frac{1}{2}t)})_{ll'}^j &= 2 \sqrt{\frac{2}{d_j}} \sum_{m=-j}^j \sum_{\substack{n_r+n_s+n_t=m \\ n'_r+n'_s+n'_t=m}} C_{n_r, n_s, n_t, n_s}^{\frac{1}{2} \frac{1}{2} l} C_{n_r+n_s, n_t, m}^{l \frac{1}{2} j} \\ &\quad C_{n'_r, n'_s, n'_r+n'_s}^{\frac{1}{2} \frac{1}{2} l'} C_{n'_r+n'_s, n'_t, m}^{l' \frac{1}{2} j} D_{n_r n'_r}^{\frac{1}{2}}(a_r) D_{n_s n'_s}^{\frac{1}{2}}(a_s) D_{n_t n'_t}^{\frac{1}{2}}(a_t). \end{aligned}$$

Using the orthogonality relation (4.17), we thus obtain

$$\begin{aligned} \langle (\chi_{(\frac{1}{2}r, \frac{1}{2}s, \frac{1}{2}t)})_{ll'}^j | p_{rst}^T \rangle &= \frac{1}{2\sqrt{2d_j}} \sum_{m=-j}^j \sum_{n_r+n_s+n_t=m} C_{n_r, n_s, n_t, n_s}^{\frac{1}{2} \frac{1}{2} l} C_{n_r+n_s, n_t, m}^{l \frac{1}{2} j} \\ &\quad \left(C_{n_t, n_r, n_r+n_t}^{\frac{1}{2} \frac{1}{2} l'} C_{n_r+n_t, n_s, m}^{l' \frac{1}{2} j} - C_{n_s, n_t, n_s+n_t}^{\frac{1}{2} \frac{1}{2} l'} C_{n_s+n_t, n_r, m}^{l' \frac{1}{2} j} \right). \end{aligned}$$

Now, (4.13) follows by computing the right-hand side for the values of j, l , and l' given above. \square

D. The costratification

According to Theorem 4.8 and Corollary 3.5, the subspaces \mathcal{V}_T and \mathcal{H}_T associated with the stratum \mathcal{P}_T are given by

$$\mathcal{V}_T = \sum_{1 \leq r < s \leq N} \text{im}(\hat{p}_{rs}^T) + \sum_{1 \leq r < s < t \leq N} \text{im}(\hat{p}_{rst}^T) \quad (4.18)$$

and

$$\mathcal{H}_T = \bigcap_{1 \leq r < s \leq N} \ker(\hat{p}_{rs}^T)^\dagger \cap \bigcap_{1 \leq r < s < t \leq N} \ker(\hat{p}_{rst}^T)^\dagger, \quad (4.19)$$

where the adjoint is taken with respect to the L^2 -scalar product $\langle \cdot, \cdot \rangle$ defined by the measure ν_{\hbar} given by (3.3).

To derive \mathcal{V}_T from (4.18) and \mathcal{H}_T from (4.19) explicitly, we simplify the notation by collecting the data $\underline{j}, j, \underline{l}$, and \underline{l}' labeling the basis functions in a multi-index

$$I := (\underline{j}; j; \underline{l}; \underline{l}').$$

Let \mathcal{I} denote the totality of all these multi-indices. According to Proposition 4.3, the structure constants of multiplication, defined by

$$\chi_{I_1}^{\mathbb{C}} \cdot \chi_{I_2}^{\mathbb{C}} = \sum_{I \in \mathcal{I}} C_{I_1 I_2}^I \chi_I^{\mathbb{C}}, \quad (4.20)$$

are given by

$$C_{I_1 I_2}^I = \sqrt{\frac{d_{\underline{j}_1} d_{\underline{j}_2} d_{\underline{j}}}{d_{\underline{j}_1} d_{\underline{j}_2} d_{\underline{j}}}} U_{\underline{j}_1, \underline{j}_2}(\underline{j}, \underline{l}; \underline{l}_1, \underline{l}_2) U_{\underline{j}_1, \underline{j}_2}(\underline{j}, \underline{l}'; \underline{l}'_1, \underline{l}'_2), \quad (4.21)$$

where $I_i = (\underline{j}_i; j_i; \underline{l}_i; \underline{l}'_i)$ and $I = (\underline{j}; j; \underline{l}; \underline{l}')$. According to (4.18), the subspace \mathcal{V}_T is spanned by the functions $p_{rs}^T \chi_I^{\mathbb{C}}$ with $1 \leq r < s \leq N$, $I \in \mathcal{I}$, and the functions $p_{rst}^T \chi_I^{\mathbb{C}}$ with $1 \leq r < s \leq N$, $I \in \mathcal{I}$. We expand

$$p_{rs}^T = \sum_{K \in \mathcal{I}} (p_{rs}^T)^K \chi_K^{\mathbb{C}}, \quad p_{rst}^T = \sum_{K \in \mathcal{I}} (p_{rst}^T)^K \chi_K^{\mathbb{C}}, \quad (4.22)$$

where the coefficients $(p_{rs}^T)^K$ and $(p_{rst}^T)^K$ are given by Lemma 4.9. Then,

$$\begin{aligned} p_{rs}^T \chi_I^{\mathbb{C}} &= \sum_{J \in \mathcal{I}} A_I^J(r, s) \chi_J^{\mathbb{C}}, & A_I^J(r, s) &:= \sum_{K \in \mathcal{I}} (p_{rs}^T)^K C_{KI}^J, \\ p_{rst}^T \chi_I^{\mathbb{C}} &= \sum_{J \in \mathcal{I}} B_I^J(r, s, t) \chi_J^{\mathbb{C}}, & B_I^J(r, s, t) &:= \sum_{K \in \mathcal{I}} (p_{rst}^T)^K C_{KI}^J. \end{aligned}$$

Thus, \mathcal{V}_T is spanned by the functions

$$\sum_{J \in \mathcal{I}} A_I^J(r, s) \chi_J^{\mathbb{C}}, \quad r < s, \quad \sum_{J \in \mathcal{I}} B_I^J(r, s, t) \chi_J^{\mathbb{C}}, \quad r < s < t, \quad I \in \mathcal{I}. \quad (4.23)$$

It remains to determine the coefficients $A_I^J(r, s)$ and $B_I^J(r, s, t)$. Recall the notation $(j_1 r_1, j_2 r_2, \dots)$ for a sequence of spins having entries j_1 at place r_1, j_2 at place r_2 , etc., and 0 elsewhere. In addition, we introduce the notation $(j_1 |_{r_1}^{s_1}, j_2 |_{r_2}^{s_2}, \dots)$ for a sequence of spins having entries j_1 at places r_1, \dots, s_1, j_2 at places r_2, \dots, s_2 , etc., and 0 elsewhere. From Lemma 4.9, we obtain

Theorem 4.10. *The vanishing subspace \mathcal{V}_T is spanned by the functions (4.23), with the coefficients $A_I^J(r, s)$ and $B_I^J(r, s, t)$ given by*

$$\begin{aligned} A_I^J(r, s) &= C_{I, ((1r, 0s); 1; (1 |_r^N); (1 |_s^N))}^J + C_{I, ((0r, 1s); 1; (1 |_s^N); (1 |_r^N))}^J \\ &\quad + C_{I, ((1r, 1s); 0; (1 |_r^{s-1}); (1 |_s^{s-1}))}^J - \frac{2}{\sqrt{3}} C_{I, ((1r, 1s); 1; (1 |_r^N); (1 |_s^N))}^J - 3 \delta_I^J, \\ B_I^J(r, s, t) &= \frac{\sqrt{3}}{2} C_{I, ((\frac{1}{2}r, \frac{1}{2}s, \frac{1}{2}t); \frac{1}{2}; (\frac{1}{2} |_r^{s-1}, 0 |_s^{t-1}, \frac{1}{2} |_t^N); (\frac{1}{2} |_r^{s-1}, 1 |_s^{t-1}, \frac{1}{2} |_t^N))}^J \\ &\quad - \frac{\sqrt{3}}{2} C_{I, ((\frac{1}{2}r, \frac{1}{2}s, \frac{1}{2}t); \frac{1}{2}; (\frac{1}{2} |_r^{s-1}, 1 |_s^{t-1}, \frac{1}{2} |_t^N); (\frac{1}{2} |_r^{s-1}, 0 |_s^{t-1}, \frac{1}{2} |_t^N))}^J. \end{aligned}$$

By taking the orthogonal complement, we obtain

Corollary 4.11. *The subspace \mathcal{H}_T associated with the stratum \mathcal{P}_T consists of the vectors $\varphi = \varphi^J \chi_J^{\mathbb{C}}$ whose coefficients φ^J are determined by the system of linear equations*

$$\sum_{J \in \mathcal{I}} \hat{A}_I^J(r, s) \varphi^J = 0, \quad r < s, \quad \sum_{J \in \mathcal{I}} \hat{B}_I^J(r, s, t) \varphi^J = 0, \quad r < s < t, \quad I \in \mathcal{I},$$

where $\hat{A}_I^J(r, s) = A_I^J(r, s) \|\chi_J^{\mathbb{C}}\|^2$ and $\hat{B}_I^J(r, s, t) = B_I^J(r, s, t) \|\chi_J^{\mathbb{C}}\|^2$, with the norm $\|\chi_J^{\mathbb{C}}\|^2$ given by (3.15).

Remark 4.12. For given multi-indices $I_1 = (\underline{j}_1; j_1; l_1; l'_1)$ and $I_2 = (\underline{j}_2; j_2; l_2; l'_2)$, the range of $I_3 = (\underline{j}_3; j_3; l_3; l'_3)$ for which the structure constant $C_{I_1 I_2}^{I_3}$ is nonzero is given by

$$\underline{j}_3 \in \langle \underline{j}_1, \underline{j}_2 \rangle, \quad j_3 \in \langle j_1, j_2 \rangle, \quad l_3, l'_3 \in R(\underline{j}_3, j_3).$$

In particular, the range is finite. Hence, the sums in Corollary 4.11 are finite. Furthermore, for fixed I_1 and I_3 , the range of I_2 is finite, too, because \underline{j}_2 and j_2 are bounded by

$$|j_1^i - j_2^i| \leq j_3^i, \quad i = 1, \dots, N, \quad |j_1 - j_2| \leq j_3$$

and the range of the sequences $\underline{l}_2, \underline{l}'_2$ of intermediate spins is given by $R(\underline{j}_2, j_2)$. Since $C_{I_1 I_2}^{I_3} = C_{I_2 I_1}^{I_3}$, for fixed I_2 and I_3 , the range of I_1 is bounded as well.

To find the coefficients φ^J , one has to rewrite the defining equations into recurrence relations and to use the asymptotic behaviour of the norms. This will be discussed elsewhere. ♦

To summarize, the costratification for $G = \text{SU}(2)$ consists of the Hilbert subspaces

$$\mathcal{H}_V, \quad \mathcal{H}_T, \quad \mathcal{H}_Z = \mathcal{H},$$

together with their orthogonal projectors. Here, \mathcal{H}_V is given by Remark 4.7 and \mathcal{H}_T is given by Theorem 4.11. By Lemma 7 of Ref. 10, the orthogonal projector onto the vanishing subspaces of the point strata is given by

$$\mathbb{P}_V(f) := f - f(v_1 \mathbb{1}, \dots, v_N \mathbb{1}) \mathbb{1},$$

where $\mathbb{1}$ denotes the constant function on $G_{\mathbb{C}}^N$ with value one. Thus, the projector onto \mathcal{H}_V is $\text{id}_{\mathcal{H}} - \mathbb{P}_V$. There are various approaches to the construction of the orthogonal projector for the T -stratum. One of them consists in applying the Schmidt orthogonalization procedure to the family (4.23). This will be studied in future work.

E. The case $N = 2$

To illustrate the general result, let us discuss the case $N = 2$. Here, for given $\underline{j} = (j^1, j^2)$ and $j \in \langle \underline{j} \rangle \equiv \langle j^1, j^2 \rangle$, the set $R(\underline{j}, j)$ consists of the single sequence $\underline{l} = (j)$. Hence, the labels \underline{l} and \underline{l}' are redundant and the basis functions may be denoted by $(\chi_{(j^1, j^2)}^{\mathbb{C}})^j$. Moreover, the isotypical components of $H_{\underline{j}} = H_{j^1} \otimes H_{j^2}$ are irreducible and the endomorphisms $A_{\underline{l}, \underline{l}'}^{j, j}$ boil down to orthogonal projectors $A^{(j^1, j^2), j}$. Thus, the basis functions are given by

$$(\chi_{(j^1, j^2)}^{\mathbb{C}})^j(a_1, a_2) = \sqrt{\frac{d_{j^1} d_{j^2}}{d_j}} \text{tr} \left(A^{(j^1, j^2), j} \circ (D^{j^1}(a_1) \otimes D^{j^2}(a_2)) \right).$$

By (3.15),

$$\|(\chi_{(j^1, j^2)}^{\mathbb{C}})^j\|^2 = (\hbar\pi)^3 e^{\hbar((2j^1+1)^2 + (2j^2+1)^2)}. \quad (4.24)$$

According to Lemma 4.4, the recoupling coefficients U are given by

$$U_{\underline{j}_1, \underline{j}_2}(j, j; j_1, j_2) = \begin{pmatrix} j_1^1 & j_2^1 & j^1 \\ j_1^2 & j_2^2 & j^2 \\ j_1 & j_2 & j \end{pmatrix}.$$

Hence, by Proposition 4.3, the multiplication law reads

$$(\chi_{\underline{j}_1}^{\mathbb{C}})^{j_1} \cdot (\chi_{\underline{j}_2}^{\mathbb{C}})^{j_2} = \sum_{\substack{j^1 \in \langle j_1^1, j_2^1 \rangle \\ j^2 \in \langle j_1^2, j_2^2 \rangle}} \sum_{j \in \langle j^1, j^2 \rangle} \begin{bmatrix} j_1^1 & j_2^1 & j^1 \\ j_1^2 & j_2^2 & j^2 \\ j_1 & j_2 & j \end{bmatrix} (\chi_{\underline{j}}^{\mathbb{C}})^j, \quad (4.25)$$

where

$$\begin{bmatrix} j_1^1 & j_2^1 & j^1 \\ j_1^2 & j_2^2 & j^2 \\ j_1 & j_2 & j \end{bmatrix} = \sqrt{\frac{d_{j_1^1} d_{j_1^2} d_{j_2^1} d_{j_2^2} d_j}{d_{j_1} d_{j_2} d_{j^1} d_{j^2}}} \begin{pmatrix} j_1^1 & j_2^1 & j^1 \\ j_1^2 & j_2^2 & j^2 \\ j_1 & j_2 & j \end{pmatrix}^2$$

is the structure constant $C_{(\underline{j}_1, j_1), (\underline{j}_2, j_2)}^{(j, j)}$. In terms of Wigner's $9j$ symbols,

$$\begin{bmatrix} j_1 & j_2 & j_3 \\ j_4 & j_5 & j_6 \\ j_7 & j_8 & j_9 \end{bmatrix} = \sqrt{\prod_{i=1}^9 d_{j_i}} \left\{ \begin{matrix} j_1 & j_2 & j_3 \\ j_4 & j_5 & j_6 \\ j_7 & j_8 & j_9 \end{matrix} \right\}^2.$$

By the analysis of Subsection IV B, we have the following orbit types and strata.

- (G) The stabilizer is $G = \text{SU}(2)$. The corresponding subset of \mathcal{P} consists of (the trivial orbits of) the points

$$((a, b), (A, B)) = (\pm 1, \pm 1), (0, 0).$$

Hence, this subset decomposes into the four strata $\mathcal{P}_{\pm\pm}$.

(T) The stabilizer is a torus. The corresponding stratum \mathcal{P}_T consists of the orbits of the points $((a_1, a_2), (A_1, A_2)) \neq ((\pm 1, \pm 1), (0, 0))$ for which a_1, a_2, A_1, A_2 commute pairwise.

(Z) The stabilizer is the center of $SU(2)$. The corresponding stratum \mathcal{P}_0 consists of the orbits of the points $((a_1, a_2), (A_1, A_2))$ for which a_1, a_2, A_1, A_2 do not commute pairwise.

To these strata, there correspond the following closed subspaces of \mathcal{H} .

(Z) As in the general case, the subspace \mathcal{H}_0 associated with the principal stratum \mathcal{P}_0 coincides with \mathcal{H} .

(G) The subspaces $\mathcal{H}_{\pm\pm}$ associated with the strata $\mathcal{P}_{\pm\pm}$ can be constructed as outlined in Remark 4.7. Since $\mathcal{P}_{\pm\pm}$ corresponds to $(\pm 1, \pm 1) \in G_{\mathbb{C}} \times G_{\mathbb{C}}$, the subspace $\mathcal{V}_{\pm\pm}$ is spanned by the functions

$$(\chi_{(j^1, j^2)}^{\mathbb{C}})^j - (\chi_{(j^1, j^2)}^{\mathbb{C}})^j (\pm 1, \pm 1)$$

and the subspace $\mathcal{H}_{\pm\pm}$ is spanned by the single vector

$$\psi_{\pm\pm} = \frac{1}{N_{\pm\pm}} \sum_{j^1, j^2} \sum_{j \in \langle j^1, j^2 \rangle} \overline{(\chi_{(j^1, j^2)}^{\mathbb{C}})^j (\pm 1, \pm 1)} (\chi_{(j^1, j^2)}^{\mathbb{C}})^j.$$

(T) The subspace \mathcal{H}_T associated with the stratum \mathcal{P}_T is defined by the single function

$$p^T(a_1, a_2) = \text{tr}([a_1, a_2]^2).$$

By Lemma 4.9,

$$p^T = (\chi_{(1,0)}^{\mathbb{C}})^1 + (\chi_{(0,1)}^{\mathbb{C}})^1 + (\chi_{(1,1)}^{\mathbb{C}})^0 - \frac{2}{\sqrt{3}} (\chi_{(1,1)}^{\mathbb{C}})^1 - 3. \quad (4.26)$$

As a consequence, the vanishing subspace \mathcal{V}_T is spanned by the vectors

$$\sum_{J \in \mathcal{I}} A_I^J \chi_J^{\mathbb{C}}, \quad I \in \mathcal{I},$$

where according to Theorem 4.10, the coefficients A_I^J are given by

$$A_I^J = \begin{bmatrix} 1 & i^1 & j^1 \\ 0 & i^2 & j^2 \\ 1 & i & j \end{bmatrix} + \begin{bmatrix} 0 & i^1 & j^1 \\ 1 & i^2 & j^2 \\ 1 & i & j \end{bmatrix} + \begin{bmatrix} 1 & i^1 & j^1 \\ 1 & i^2 & j^2 \\ 0 & i & j \end{bmatrix} - \frac{2}{\sqrt{3}} \begin{bmatrix} 1 & i^1 & j^1 \\ 1 & i^2 & j^2 \\ 1 & i & j \end{bmatrix} - 3\delta_I^J$$

with $L = (i; i)$ and $J = (j; j)$. Finally, Corollary 4.11 implies that \mathcal{H}_T consists of the functions $\varphi = \varphi^J \chi_J^{\mathbb{C}}$ whose coefficients φ^J are determined by the system of linear equations

$$\sum_{J \in \mathcal{I}} \hat{A}_I^J \varphi^J = 0, \quad I \in \mathcal{I},$$

where $\hat{A}_I^J = A_I^J \|\chi_J^{\mathbb{C}}\|^2$ and $\|\chi_J^{\mathbb{C}}\|^2$ is given by (4.24).

F. The eigenvalue problem for the Hamiltonian

Recall the classical Hamiltonian

$$H(a, E) = \frac{g^2}{2\delta} \sum_{\ell \in \Lambda^1} \|E(\ell)\|^2 - \frac{1}{g^2\delta} \sum_{p \in \Lambda^2} (\text{tr } a(p) + \overline{\text{tr } a(p)}),$$

given by (2.2). Here, $a(p) = a(\ell_1)a(\ell_2)a(\ell_3)a(\ell_4)$, where the links ℓ_1, \dots, ℓ_4 , in this order, form the boundary of p and are endowed with the boundary orientation. The quantum Hamiltonian, obtained via canonical quantization in the tree gauge, is called the Kogut-Susskind Hamiltonian (more precisely, its pure gauge part),

$$H = \frac{g^2}{2\delta} \mathfrak{C} - \frac{1}{g^2\delta} \mathfrak{W}. \quad (4.27)$$

Here,

$$\mathfrak{C} := \sum_{\ell \in \Lambda^1} E_{ij}(\ell) E_{ji}(\ell)$$

is the Casimir operator (negative of the group Laplacian) of $SU(2)^N$ and

$$\mathfrak{W} := \sum_{p \in \Lambda^2} (W(p) + W(p)^*),$$

where $W(p)$ is the quantum counterpart (multiplication operator on \mathcal{H}) of $\text{tr } a(p)$, called the Wilson loop operator. For details, see Refs. 26, 14, and 27. Recall that the representative functions of spin j on $SU(2)$ are eigenfunctions of the Casimir operator of $SU(2)$ corresponding to the eigenvalue

$$\epsilon_j = 4j(j+1),$$

see Refs. 8 and 18. It follows that the invariant representative functions $(\chi_{\underline{j}})_{\underline{l}, \underline{l}'}$ are eigenfunctions of \mathfrak{C} corresponding to the eigenvalues

$$\epsilon_{\underline{j}} = \epsilon_{j_1} + \dots + \epsilon_{j_N}. \quad (4.28)$$

Let us analyze \mathfrak{W} . For this purpose, for our regular cubic lattice, we define a standard tree as follows. By a line, we mean a maximal straight line consisting of lattice links. First, choose a lattice site x_0 and a line L_1 through x_0 . Next, choose a second line L_2 through x_0 perpendicular to L_1 and add all lines parallel to L_2 in the plane spanned by L_1 and L_2 . Finally, add all lines perpendicular to that plane. Let B be such a standard tree. Since $a(\ell) = 1$ for every $\ell \in B^1$ and since there are no plaquettes having 3 off-tree links, we can decompose \mathfrak{W} into three sums. It is easy to check that there exists an orientation and a numbering of the off-tree links such that for every plaquette with four off-tree links (all of these plaquettes are parallel to the plane spanned by the lines L_1 and L_2), the boundary links are numbered and oriented consistently, meaning that for one of the two possible orientations of the plaquette, they carry the induced boundary orientation and that their numbers increase in that direction. Then,

$$\begin{aligned} \mathfrak{W} = & \sum_{\{p: p \cap B = \emptyset\}} \text{tr}(a_{r_p} a_{s_p} a_{t_p} a_{u_p}) + \overline{\text{tr}(a_{r_p} a_{s_p} a_{t_p} a_{u_p})} \\ & + \sum_{\{p: |p \cap B| = 2\}} \text{tr}(a_{r_p} a_{s_p}) + \overline{\text{tr}(a_{r_p} a_{s_p})} \\ & + \sum_{\{p: |p \cap B| = 3\}} \text{tr}(a_{r_p}) + \overline{\text{tr}(a_{r_p})}. \end{aligned}$$

To find the matrix elements of H with respect to the basis functions $\{(\chi_{\underline{j}})_{\underline{l}, \underline{l}'}^{\underline{j}}\}$, we have to find the corresponding expansion of \mathfrak{W} . The sequences \underline{j} occurring here will have at most four nonzero entries, so we can use the notation introduced in Sec. IV C, given by writing $\underline{j} = (j_1 r_1, \dots, j_k r_k)$ if \underline{j} has entries 0, except for j_i at the places r_i , $i = 1, \dots, k$. The function $T_r(\underline{a}) = \text{tr}(a_r)$ coincides with the basis function $(\chi_{\underline{j}})_{\underline{l}, \underline{l}'}^{\underline{j}}$ with $\underline{j} = (\frac{1}{2} r)$. Omitting the irrelevant indices $\underline{j}, \underline{l}, \underline{l}'$, we thus have

$$T_r = \chi_{(\frac{1}{2} r)}.$$

The function $T_{rs}(\underline{a}) := \text{tr}(a_r a_s)$ is a linear combination of the basis functions $(\chi_{\underline{j}})_{\underline{l}, \underline{l}'}^{\underline{j}}$ with $\underline{j} = (\frac{1}{2} r, \frac{1}{2} s)$. As in Sec. IV C, we may omit the irrelevant labels $\underline{l}, \underline{l}'$. Using (4.16), the orthogonality relation for the matrix entry functions given by (4.17) and the normalization condition $\sum_{m_1, m_2 = \pm \frac{1}{2}} (C_{m_1 m_2 m}^{\frac{1}{2} \frac{1}{2} j})^2 = 1$, we obtain

$$\langle (\chi_{(\frac{1}{2} r, \frac{1}{2} s)})^{\underline{j}} | T_{rs} \rangle = (-1)^{1-j} \frac{\sqrt{d_j}}{2}.$$

Thus,

$$T_{rs} = \frac{\sqrt{3}}{2} (\chi_{(\frac{1}{2} r, \frac{1}{2} s)})^1 - \frac{1}{2} (\chi_{(\frac{1}{2} r, \frac{1}{2} s)})^0$$

Finally, the function $T_{rstu}(\underline{a}) := \text{tr}(a_r a_s a_t a_u)$ is a linear combination of the basis functions $(\chi_j)_{ll'}^j$ with $j = (\frac{1}{2}r, \frac{1}{2}s, \frac{1}{2}t, \frac{1}{2}u)$. Here, the sequences $\underline{l} \in R(j, j)$ have entries $l^1 = \dots = l^{r-1} = 0$, $l^r = \dots = l^{s-1} = \frac{1}{2}$, $l^s = \dots = l^{t-1} = l$, $l^t = \dots = l^{u-1} = k$, and $l^u = \dots = l^N = j$, where $l = 0, 1$ and $k \in \langle l, \frac{1}{2} \rangle$ so that $j \in \langle k, \frac{1}{2} \rangle$. That is, they are labeled by two intermediate spins l, k so that in our notation we may replace the labels \underline{l} and \underline{l}' by (l, k) and (l', k') , respectively. Expressing these basis functions in terms of matrix entry functions according to Proposition 4.2 and using once again the orthogonality relation (4.17), we obtain

$$\begin{aligned} \langle (\chi_{(\frac{1}{2}r, \dots, \frac{1}{2}u)})_{(l,k),(l',k')}^j | T_{rstu} \rangle &= \frac{1}{4\sqrt{d_j}} \sum_{m=-j}^j \sum_{m_r + \dots + m_u = m} \\ &\quad C_{m_r, m_s, m_t, m_u}^{l, \frac{1}{2}l} C_{m_r + m_s, m_t, m_r + m_s + m_t}^{l, \frac{1}{2}k} C_{m_r + m_s + m_t, m_u, m}^{k, \frac{1}{2}j} \\ &\quad C_{m_u, m_r, m_r + m_u}^{\frac{1}{2}l, \frac{1}{2}l'} C_{m_r + m_u, m_s, m_r + m_s + m_u}^{l', \frac{1}{2}k'} C_{m_r + m_s + m_u, m_t, m}^{k', \frac{1}{2}j}. \end{aligned}$$

Evaluation yields

$$\begin{aligned} T_{rstu} &= \frac{1}{8} (\chi_{(\frac{1}{2}r, \dots, \frac{1}{2}u)})_{(0\frac{1}{2})(0\frac{1}{2})}^0 - \frac{\sqrt{3}}{8} (\chi_{(\frac{1}{2}r, \dots, \frac{1}{2}u)})_{(1\frac{1}{2})(0\frac{1}{2})}^0 - \frac{\sqrt{3}}{8} (\chi_{(\frac{1}{2}r, \dots, \frac{1}{2}u)})_{(0\frac{1}{2})(1\frac{1}{2})}^0 \\ &\quad - \frac{1}{8} (\chi_{(\frac{1}{2}r, \dots, \frac{1}{2}u)})_{(1\frac{1}{2})(1\frac{1}{2})}^0 - \frac{\sqrt{3}}{8} (\chi_{(\frac{1}{2}r, \dots, \frac{1}{2}u)})_{(0\frac{1}{2})(0\frac{1}{2})}^1 - \frac{1}{8} (\chi_{(\frac{1}{2}r, \dots, \frac{1}{2}u)})_{(1\frac{1}{2})(0\frac{1}{2})}^1 \\ &\quad - \frac{1}{\sqrt{8}} (\chi_{(\frac{1}{2}r, \dots, \frac{1}{2}u)})_{(1\frac{3}{2})(0\frac{1}{2})}^1 + \frac{3}{8} (\chi_{(\frac{1}{2}r, \dots, \frac{1}{2}u)})_{(0\frac{1}{2})(1\frac{1}{2})}^1 - \frac{1}{8\sqrt{3}} (\chi_{(\frac{1}{2}r, \dots, \frac{1}{2}u)})_{(1\frac{1}{2})(1\frac{1}{2})}^1 \\ &\quad - \frac{1}{2\sqrt{6}} (\chi_{(\frac{1}{2}r, \dots, \frac{1}{2}u)})_{(1\frac{3}{2})(1\frac{1}{2})}^1 + \frac{1}{\sqrt{6}} (\chi_{(\frac{1}{2}r, \dots, \frac{1}{2}u)})_{(1\frac{1}{2})(1\frac{3}{2})}^1 - \frac{1}{4\sqrt{3}} (\chi_{(\frac{1}{2}r, \dots, \frac{1}{2}u)})_{(1\frac{3}{2})(1\frac{3}{2})}^1 \\ &\quad + \frac{\sqrt{5}}{4} (\chi_{(\frac{1}{2}r, \dots, \frac{1}{2}u)})_{(1\frac{3}{2})(1\frac{3}{2})}^2. \end{aligned}$$

Now, consider the eigenvalue problem for H . Expanding

$$\psi = \sum_J \psi^J \chi_J, \quad \mathfrak{W} = \sum_I W^I (\chi_I + \overline{\chi_I}),$$

and using (4.20), as well as the fact that $\langle \chi_K | \overline{\chi_I} \chi_J \rangle = \langle \chi_I \chi_K | \chi_J \rangle = C_{IK}^J$ implies

$$\overline{\chi_I} \chi_J = C_{IK}^J \chi_K,$$

we can write the eigenvalue equation in the form

$$\sum_{J \in \mathcal{I}} \left\{ \left(\frac{g^2}{2\delta} \epsilon_J - \mathcal{E} \right) \delta_J^K - \frac{1}{g^2 \delta} \sum_{I \in \mathcal{I}} W^I (C_{IJ}^K + C_{IK}^J) \right\} \psi^J = 0, \quad (4.29)$$

for all $K \in \mathcal{I}$. Here, we have written ϵ_J for the eigenvalue of the Casimir operator \mathfrak{C} corresponding to the eigenfunction χ_J , given by (4.28). Thus, we are left with a homogeneous system of linear equations for the eigenfunction coefficients ψ^J . The eigenvalues \mathcal{E} are determined by the requirement that the determinant of this system must vanish. Note that the sum over I in (4.29) is finite because there are only finitely many nonvanishing W^I . Moreover, by Remark 4.12, also the sum over J is finite for every fixed K . Thus, we have reduced the eigenvalue problem for the Hamiltonian to a problem in linear algebra. Combining this with the well-known asymptotic properties of $3nj$ symbols, see Ref. 5 (Topic 9), Ref. 2, and further references therein, we obtain an algebraic setting which allows for a computer algebra supported study of the spectral properties of H . This will be done in a future work.

V. SUMMARY AND OUTLOOK

In this paper, we have constructed the Hilbert space costratification for $SU(2)$ lattice gauge theory. This work is based on the results obtained in Ref. 11, where we have implemented the

defining relations for the orbit type strata on the quantum level. Here, the main technical tool is the calculus of invariant representative functions for representations of $SU(2)$ combined with recoupling theory for angular momenta. We have already explained in the Introduction how the results of this paper fit into our long-term programme for studying the non-perturbative aspects of non-Abelian quantum gauge theories. Here, let us outline some perspectives:

1. It will be a challenge to extend our results to the case of the gauge group $SU(3)$. On the classical level, we have some preliminary results; see, e.g., the case studies in Refs. 6, 7, and 9.
2. In Ref. 10, one of us has developed another approach toward the study of costratifications for arbitrary compact Lie groups. The starting point in Ref. 10 was the observation that the vanishing subspaces corresponding to the classical strata may be viewed as intersections of one-point vanishing subspaces. The orthogonal complements of the latter were shown to be one-dimensional and, for each of these spaces, a spanning (holomorphic and square integrable) vector w_g was constructed. Finally, passing to the Hilbert space of invariant functions was accomplished by using the projection operator \mathbb{P} defined by averaging over the compact group manifold. As a result, each element of the costratification was characterized as the closure of the span of $\{\mathbb{P}w_g\}$, with g running over a complete set of representatives of the set of orbits belonging to the stratum under consideration. Moreover, for the point strata, the spanning vectors w_g turned out to be proportional to the coherent states in the sense of Hall.^{15,16} For the other strata, up until now, this approach has not led to such an explicit characterization of the corresponding elements of the costratification. It will be interesting to combine the calculus developed in this paper with the methods of Ref. 10. This will possibly lead to a characterization of the full costratification in terms of coherent like states.
3. In Subsection IV F, we have formulated the eigenvalue problem of the quantum Hamiltonian H in terms of invariant representative functions. We have shown that, in this language, it boils down to a problem in linear algebra. As already explained there, this can serve as a starting point for a study of the spectral properties of H . In particular, it should be possible to investigate the role of the coherent states addressed in the previous point; see the toy model in Ref. 22.

APPENDIX A: PROOF OF COROLLARY 3.8

For every $\lambda \in \widehat{G}$, we choose a scalar product in H_λ invariant under π_λ and an orthonormal basis $\{e_r^\lambda : r = 1, \dots, \dim(H_\lambda)\}$. For every $\underline{\lambda} \in \widehat{G}^N$, the vectors

$$e_{\underline{r}}^{\underline{\lambda}} := e_{r^1}^{\lambda^1} \otimes \cdots \otimes e_{r^N}^{\lambda^N}, \quad \underline{r} = (r^1, \dots, r^N), \quad r^i = 1, \dots, \dim(H_{\lambda^i}),$$

form an orthonormal basis in $H_{\underline{\lambda}}$ with respect to the natural scalar product in the tensor product of Hilbert spaces. Define holomorphic functions

$$f_{r,s}^\lambda : G_{\mathbb{C}} \rightarrow \mathbb{C}, \quad f_{r,s}^\lambda(a) := \sqrt{\dim(H_\lambda)} \langle e_r^\lambda | \pi_\lambda(a) e_s^\lambda \rangle,$$

where $\lambda \in \widehat{G}$, $r, s = 1, \dots, \dim(H_\lambda)$, and

$$f_{\underline{r},\underline{s}}^{\underline{\lambda}} : G_{\mathbb{C}}^N \rightarrow \mathbb{C}, \quad f_{\underline{r},\underline{s}}^{\underline{\lambda}}(a) := \sqrt{\dim(H_{\underline{\lambda}})} \langle e_{\underline{r}}^{\underline{\lambda}} | \pi_{\underline{\lambda}}(a) e_{\underline{s}}^{\underline{\lambda}} \rangle,$$

where $\underline{\lambda} \in \widehat{G}^N$ and $\underline{r}, \underline{s} \in \prod_{r=1}^N \{1, \dots, \dim(H_{\lambda^r})\}$. We have

$$f_{\underline{r},\underline{s}}^{\underline{\lambda}}(a) = \prod_{i=1}^N f_{r^i,s^i}^{\lambda^i}(a_i).$$

By this relation, the natural unitary isomorphism $\mathcal{HL}^2(G_{\mathbb{C}}^N) = (\mathcal{HL}^2(G_{\mathbb{C}}))^{\otimes N}$ and the holomorphic Peter-Weyl theorem²¹ for $G_{\mathbb{C}}$, the functions $f_{\underline{r},\underline{s}}^{\underline{\lambda}}$ form an orthogonal basis in $\mathcal{HL}^2(G_{\mathbb{C}}^N)$ and have the norms

$$\|f_{\underline{r},\underline{s}}^{\lambda}\|^2 = \prod_{i=1}^N \|f_{r^i,s^i}^{\lambda^i}\|^2 = \prod_{i=1}^N C_{\lambda^i}.$$

Using the basis vectors $e_{\underline{k}}^{\lambda}$ in $H_{\underline{\lambda}}$ to compute the trace, we find the expansion

$$(\chi_{\underline{\lambda}}^{\mathbb{C}})_{k,l}^{\lambda} = \sum_{\underline{r},\underline{s}} \langle e_{\underline{s}}^{\lambda} | A_{k,l}^{\lambda,\lambda} e_{\underline{r}}^{\lambda} | f_{\underline{r},\underline{s}}^{\lambda} \rangle. \quad (\text{A1})$$

Using this, orthogonality of the functions $f_{\underline{r},\underline{s}}^{\lambda}$, and orthonormality of the endomorphisms $A_{k,l}^{\lambda,\lambda}$, we obtain

$$\begin{aligned} \langle (\chi_{\underline{\lambda}}^{\mathbb{C}})_{k,l}^{\lambda} | (\chi_{\underline{\lambda}'}^{\mathbb{C}})_{k',l'}^{\lambda'} \rangle &= \delta_{\underline{\lambda},\underline{\lambda}'} \sum_{\underline{r},\underline{s}} \overline{\langle e_{\underline{s}}^{\lambda} | A_{k,l}^{\lambda,\lambda} e_{\underline{r}}^{\lambda} \rangle} \langle e_{\underline{s}}^{\lambda} | A_{k',l'}^{\lambda,\lambda'} e_{\underline{r}}^{\lambda} \rangle \prod_{i=1}^N C_{\lambda^i} \\ &= \delta_{\underline{\lambda},\underline{\lambda}'} \text{tr} \left((A_{k,l}^{\lambda,\lambda})^* A_{k',l'}^{\lambda,\lambda'} \right) \prod_{i=1}^N C_{\lambda^i} \\ &= \delta_{\underline{\lambda},\underline{\lambda}'} \delta_{\lambda\lambda'} \delta_{kk'} \delta_{ll'} \prod_{i=1}^N C_{\lambda^i}. \end{aligned}$$

This yields the assertion. \square

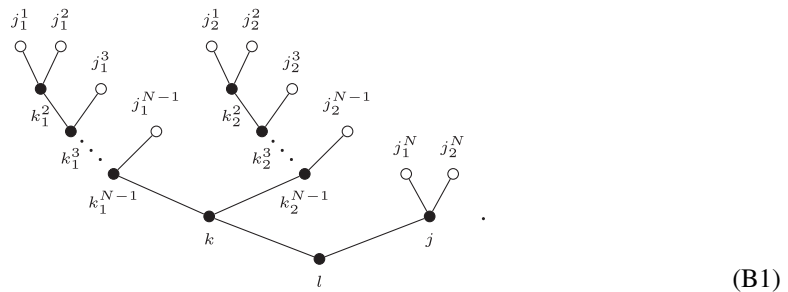
APPENDIX B: PROOF OF LEMMA 4.4

Proof. The proof is by induction over N . In the case $N = 2$, the reduction scheme (4.6) boils down to (4.10) and the reduction scheme (4.7) boils down to (4.9), where $j_1 = j_1^1, j_2 = j_2^1, j_3 = j_1^1, j_4 = j_1^2, j_5 = j_2^2, j_6 = j_2^2, j_7 = l_1^2, j_8 = l_2^2$, and $j_9 = l^2$. Hence, by definition, the scalar product of the corresponding ladder basis elements yields the recoupling coefficient

$$\begin{pmatrix} j_1^1 & j_2^1 & j_1^1 \\ j_1^2 & j_2^2 & j^2 \\ l_1^2 & l_2^2 & l^2 \end{pmatrix},$$

and by (4.8), this coincides with $U_{j_1 j_2}(j, l; l_1, l_2)$ in the case at hand. This proves the assertion for $N = 2$.

Now, let $N > 2$ be given and assume that the assertion holds for $N - 1$. In what follows, for any given sequence $\underline{x} = (x^1, \dots, x^N)$, we denote $\underline{x} := (x^1, \dots, x^{N-1})$. Recall that in $H_{j_1} \otimes H_{j_2}$, we have the orthonormal basis vectors $|j_1, j_2; l_1, l_2; l, m\rangle$ defined by the reduction scheme (4.6) and $|j_1, j_2; j, l, m\rangle$ defined by the reduction scheme (4.7). Consider the following reduction scheme:



It leads to irreducible subspaces labeled by $\underline{k}_1, \underline{k}_2, k, j, l$. Let $|\underline{k}_1, \underline{k}_2; k, j; l^N, m\rangle$ denote the elements of the corresponding orthonormal ladder basis. Using this basis to plug a unit into (4.8), we obtain

$$\begin{aligned} U_{j_1 j_2}(j, l; l_1, l_2; l^N) &= \sum_{\underline{k}_1, \underline{k}_2, k, j} \langle j_1, j_2; j, l, m | \underline{k}_1, \underline{k}_2; k, j; l^N, m \rangle \times \dots \\ &\quad \dots \times \langle \underline{k}_1, \underline{k}_2; k, j; l^N, m | j_1, j_2; l_1, l_2; l^N, m \rangle. \end{aligned} \quad (\text{B2})$$

To compute $\langle j_1, j_2; j, l, m | k_1, k_2; k, j; l^N, m \rangle$, we view $H_{j_1} \otimes H_{j_2}$ as the tensor product $(H_{j_1} \otimes H_{j_2}) \otimes (H_{j_1^N} \otimes H_{j_2^N})$ and expand both arguments with respect to appropriately chosen product bases,

$$\begin{aligned} |j_1, j_2; j, l, m\rangle &= \sum_{m_1+m_2=m} C_{m_1 m_2 m}^{l^{N-1} j^N l^N} |j_1, j_2; j, l, m_1\rangle \otimes |j_1^N, j_2^N; j^N, m_2\rangle, \\ |k_1, k_2; k, j; l^N, m\rangle &= \sum_{m_1+m_2=m} C_{m_1 m_2 m}^{k j l^N} |j_1, j_2; k_1, k_2; k, m_1\rangle \otimes |j_1^N, j_2^N; j, m_2\rangle. \end{aligned}$$

In view of (4.8), this yields

$$\langle j_1, j_2; j, l, m | k_1, k_2; k, j; l^N, m \rangle = \delta_{l^{N-1}, k} \delta_{j^N, j} U_{j_1, j_2}(\underline{j}, \underline{l}; \underline{k}_1, \underline{k}_2, k). \quad (\text{B3})$$

To compute the scalar product $\langle k_1, k_2; k, j; l^N, m | j_1, j_2; l_1, l_2; l^N, m \rangle$, we observe that the vectors $|k_1, k_2; k, j; l^N, m\rangle$ are ladder basis elements in the tensor product

$$H_{j_1, k_1} \otimes H_{j_2, k_2} \otimes H_{j_1^N} \otimes H_{j_2^N}$$

[for the notation H_{j_1, k_1} , etc., see (4.2)] defined by the reduction scheme (4.9) with $j_1 = k_1^{N-1}$, $j_2 = k_2^{N-1}$, $j_3 = k$, $j_4 = j_1^N$, $j_5 = j_2^N$, $j_6 = j$, and $j_9 = l^N$, whereas $|j_1, j_2; l_1, l_2; l^N, m\rangle$ are ladder basis elements in the tensor product

$$H_{j_1, l_1} \otimes H_{j_2, l_2} \otimes H_{j_1^N} \otimes H_{j_2^N}$$

defined by the reduction scheme (4.10) with $j_1 = l_1^{N-1}$, $j_2 = l_2^{N-1}$, $j_4 = j_1^N$, $j_5 = j_2^N$, $j_7 = l_1^N$, $j_8 = l_2^N$, and $j_9 = l^N$. Therefore,

$$\langle k_1, k_2; k, j; l^N, m | j_1, j_2; l_1, l_2; l^N, m \rangle = \delta_{k_1, l_1} \delta_{k_2, l_2} \begin{pmatrix} l_1^{N-1} & l_2^{N-1} & k \\ j_1^N & j_2^N & j \\ l_1^N & l_2^N & l^N \end{pmatrix}.$$

Plugging this and (B3) into (B2) and taking the sum, we obtain

$$U_{j_1, j_2}(\underline{j}, \underline{l}; \underline{l}_1, \underline{l}_2) = \begin{pmatrix} l_1^{N-1} & l_2^{N-1} & l^{N-1} \\ j_1^N & j_2^N & j^N \\ l_1^N & l_2^N & l^N \end{pmatrix} U_{j_1, j_2}(\underline{j}, \underline{l}; \underline{l}_1, \underline{l}_2).$$

Thus, the induction assumption implies that the assertion holds for N . \square

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